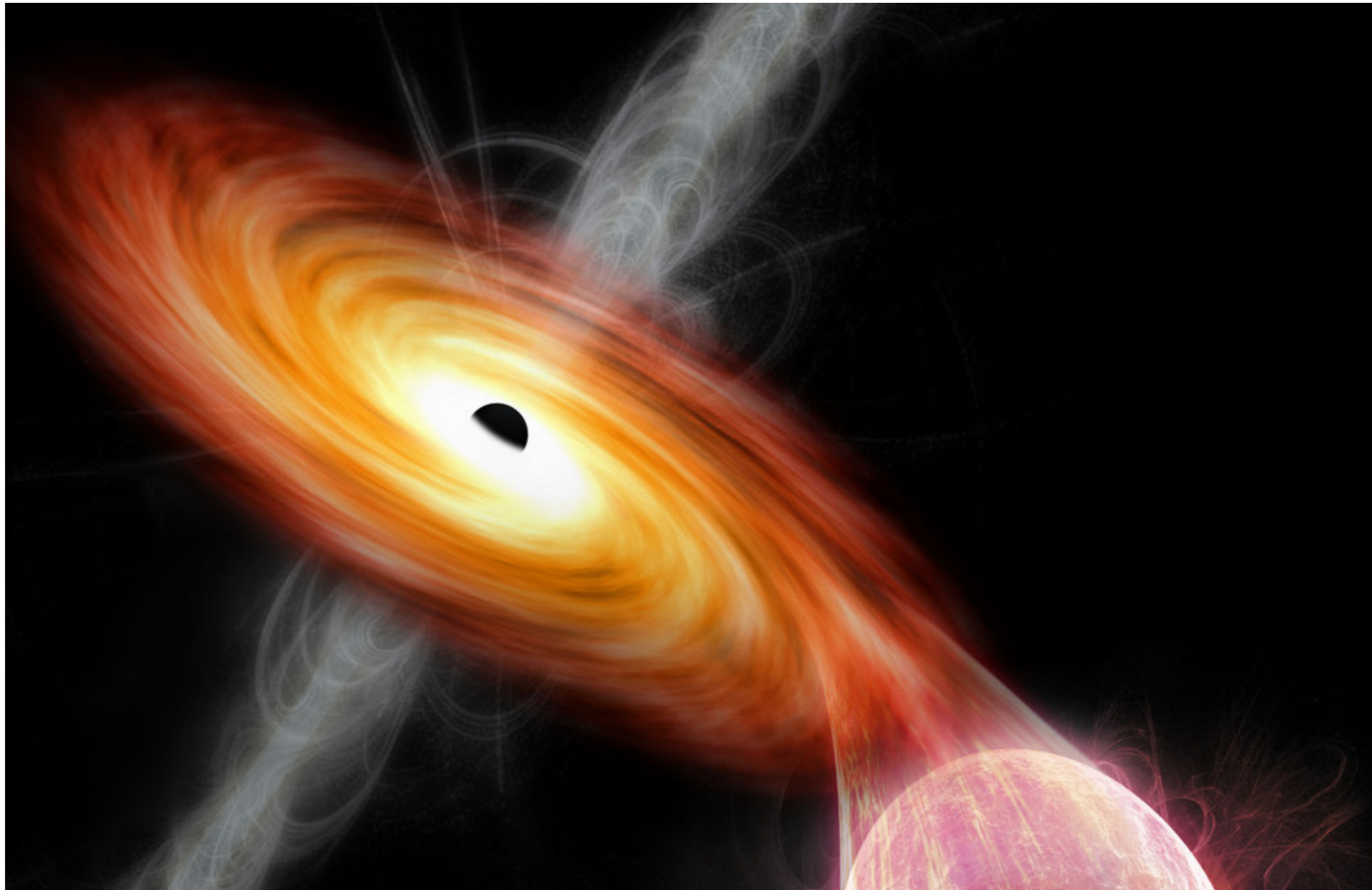


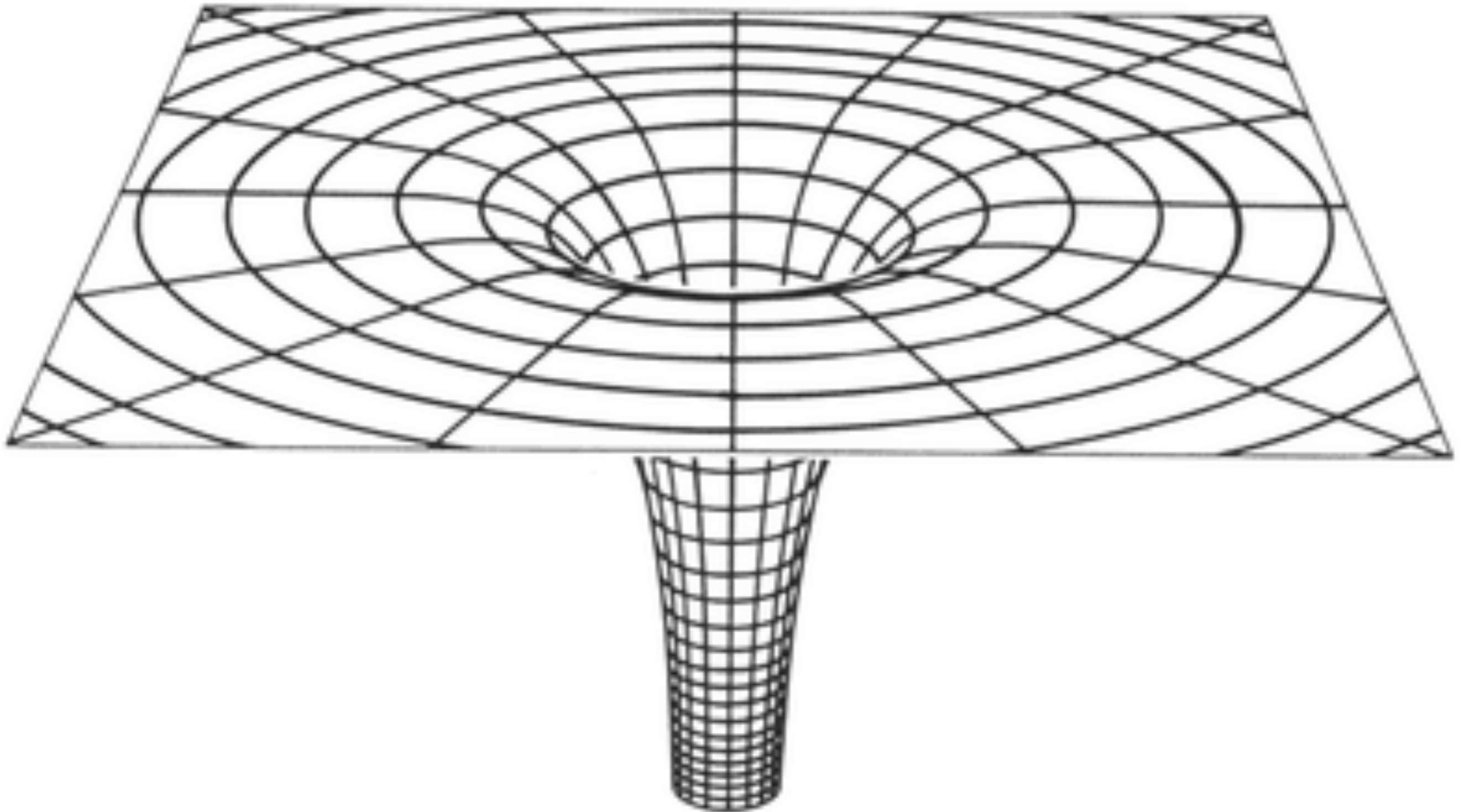
High Energy Astrophysics

Dr. Adam Ingram



Lecture 5

Black Holes



Black Holes for Babies

B is for Black Hole

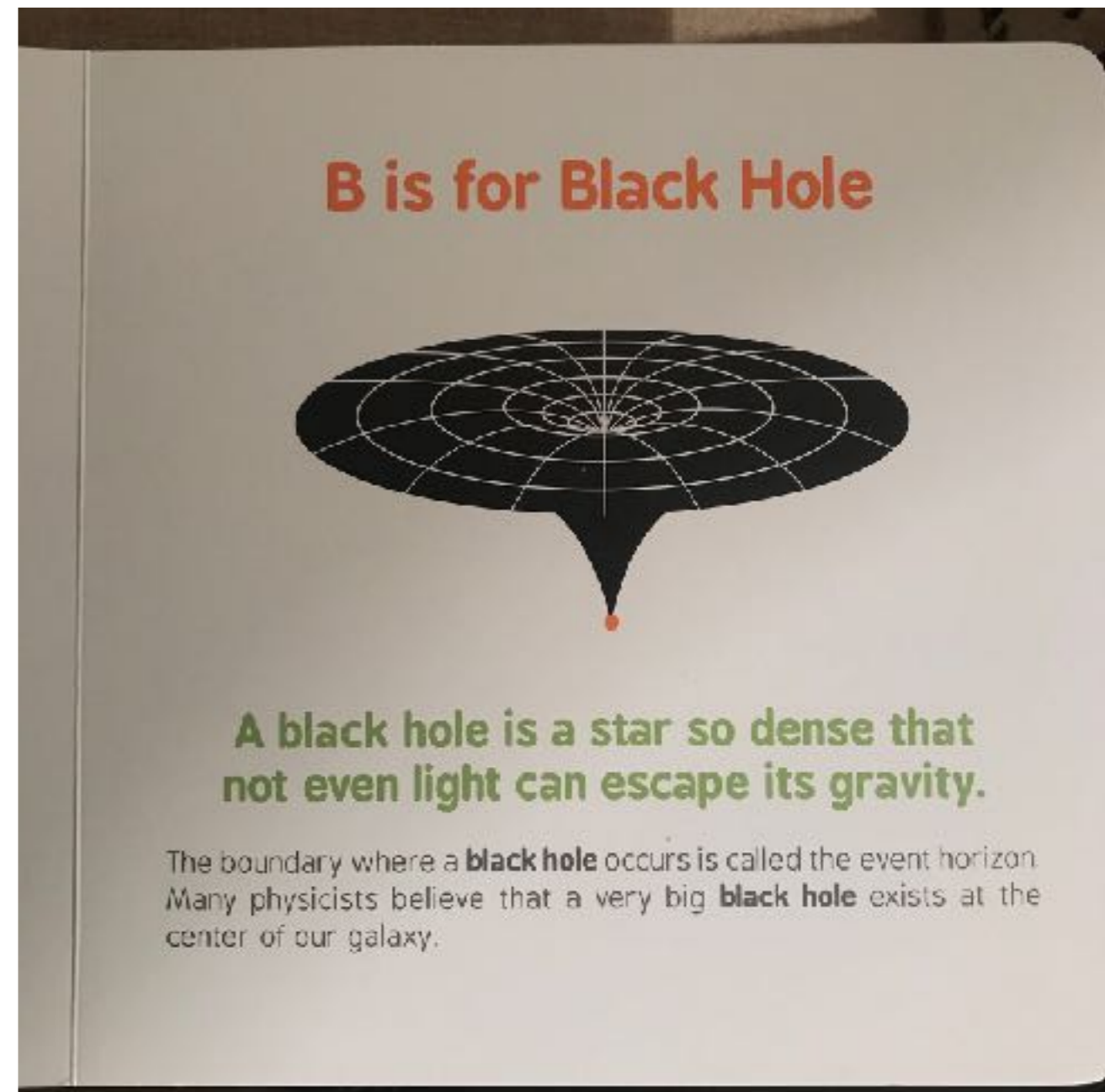


A black hole is a star so dense that not even light can escape its gravity.

The boundary where a **black hole** occurs is called the event horizon. Many physicists believe that a very big **black hole** exists at the center of our galaxy.

Black Holes for Babies

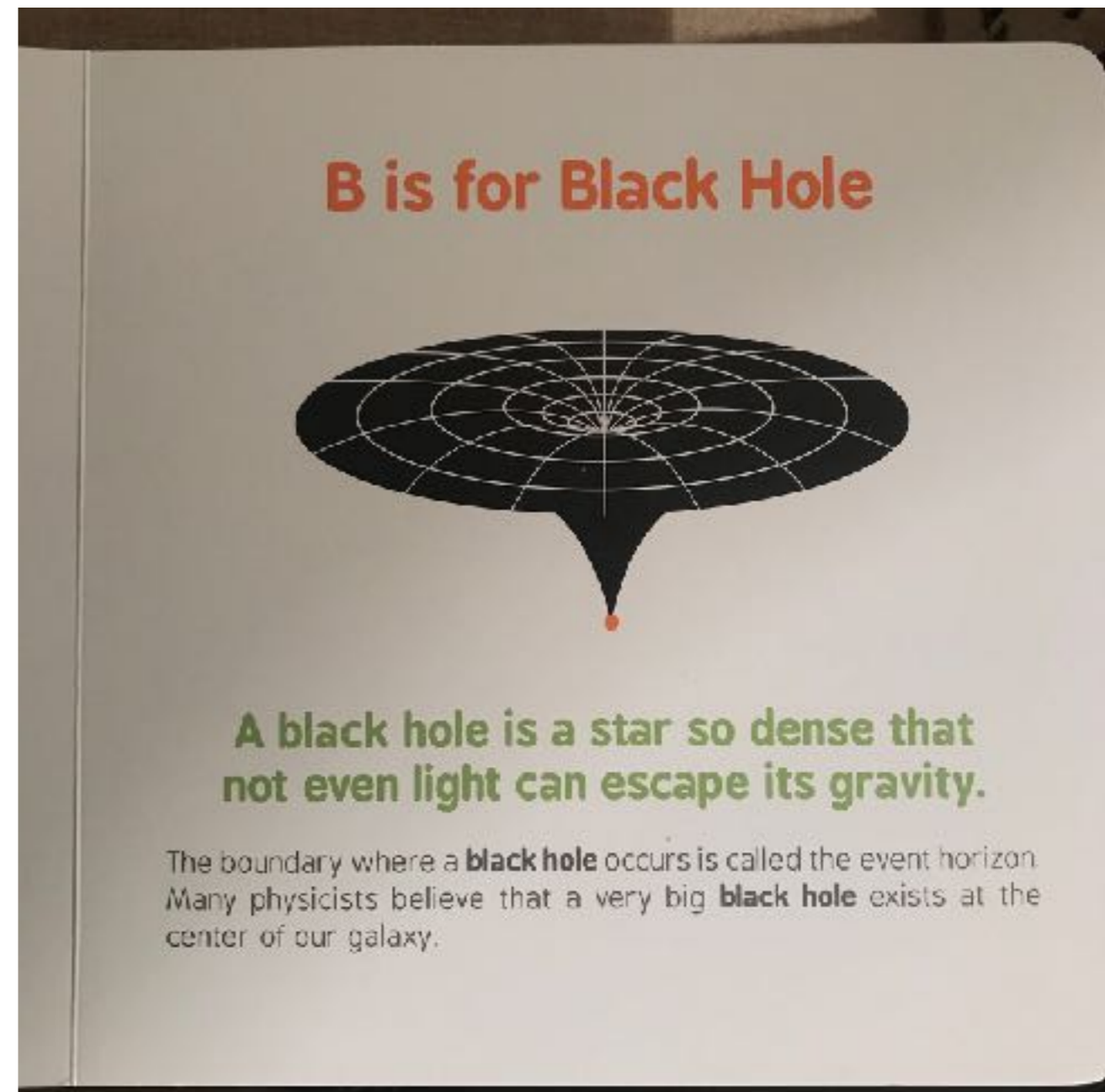
- All mass, M , in a singularity
- Event Horizon: $v_{\text{esc}} = c!$
- Newtonian approx: $v_{\text{esc}}^2 = 2GM/r \implies r_h = 2GM/c^2$
- Size scale: gravitational radius: $r_g = GM/c^2$



Black Holes for Babies

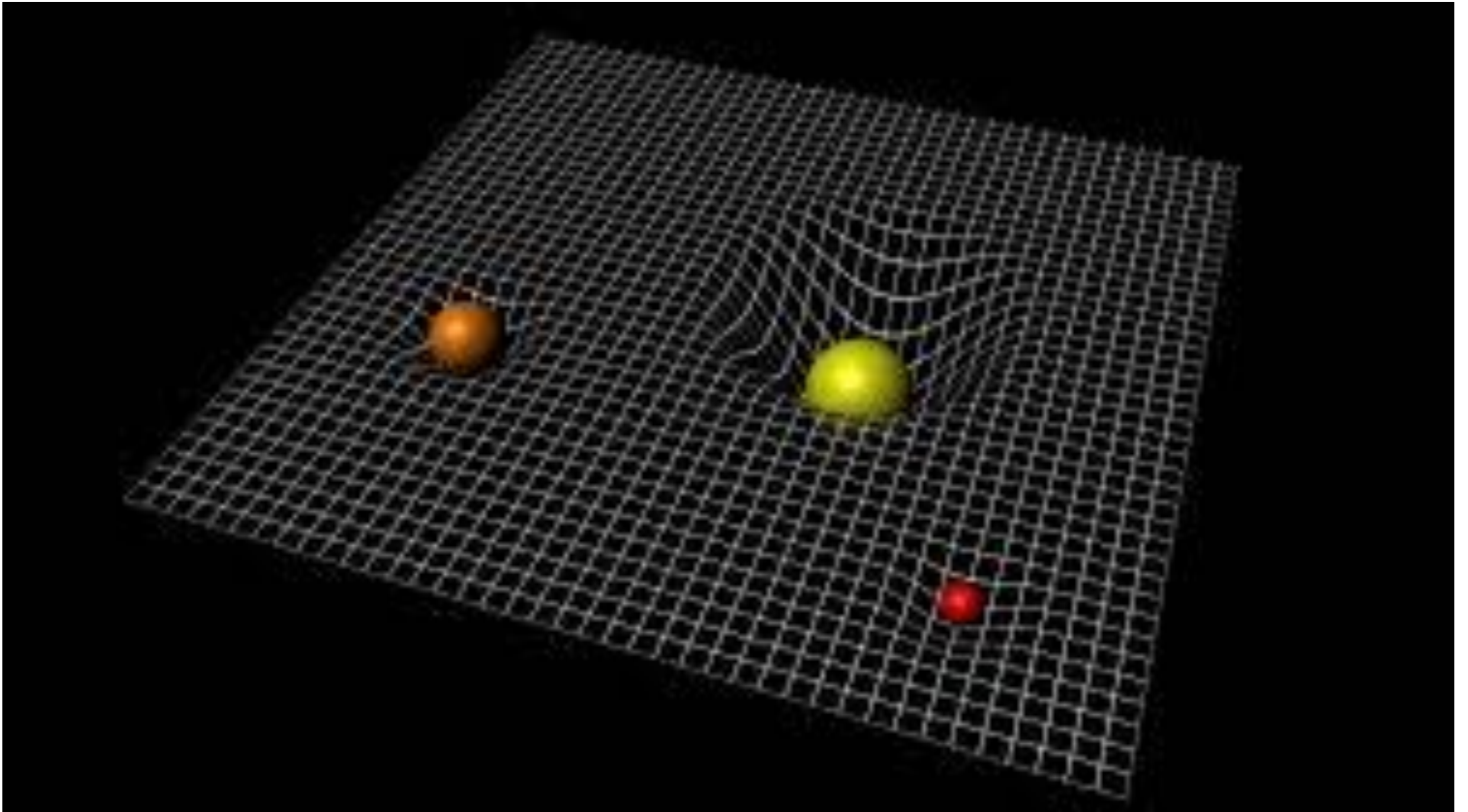
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**Symbol r_g was gyroradius;
is now gravitational radius!**



Black Holes in GR

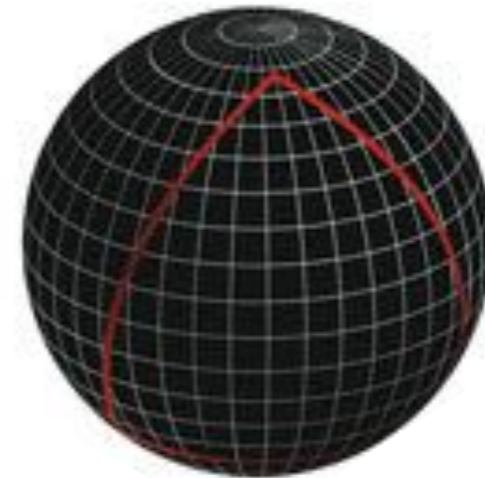
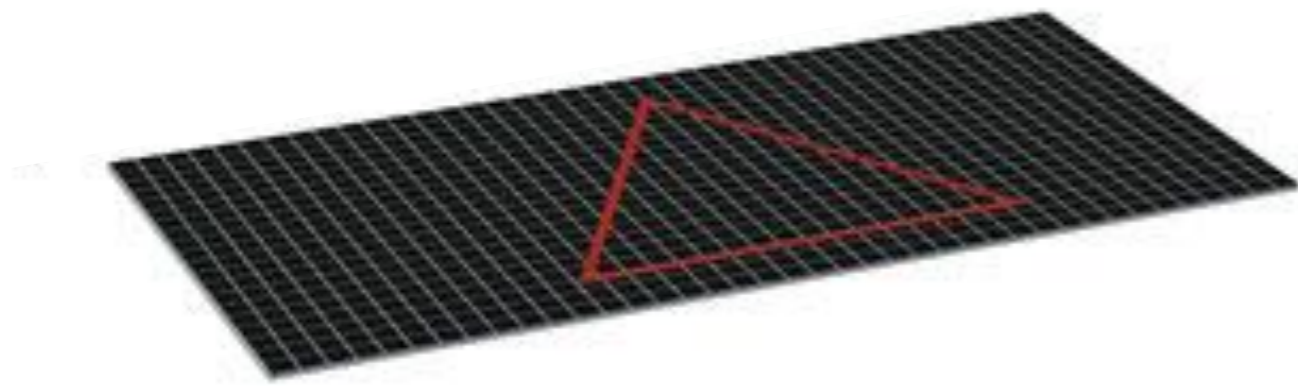
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- What does spacetime curvature mean? Well, in flat 3D space we have that the distance between two points is given by Pythagoras' theorem:

$$(d\ell)^2 = (dx)^2 + (dy)^2 + (dz)^2 = (dr)^2 + (rd\theta)^2 + (r \sin \theta d\phi)^2$$

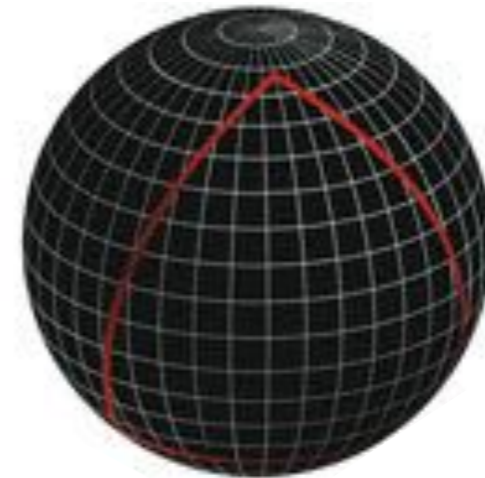
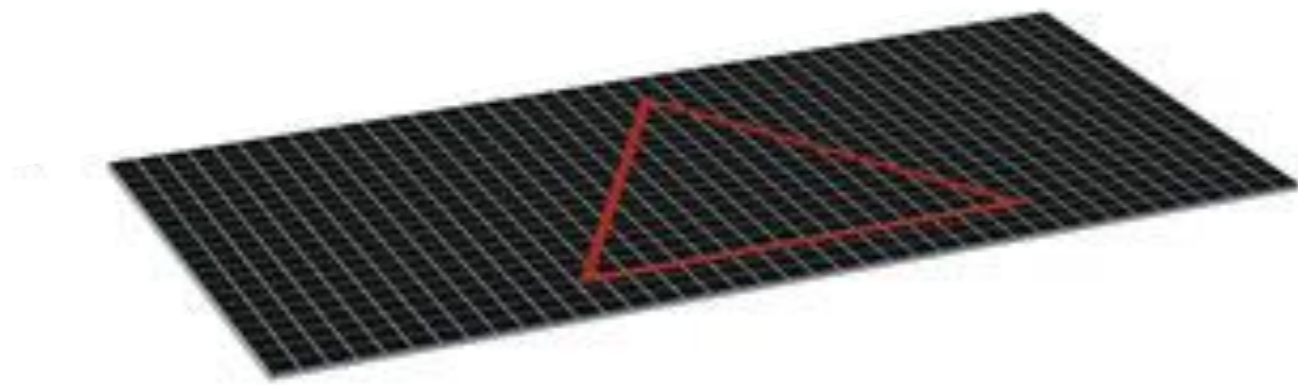


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- In **curved** 3D space, it is not given by Pythagoras.
- In SR, include time into 4D spacetime — introduce **spacetime interval**. For flat spacetime, this is:

$$(ds)^2 = - (cdt)^2 + (dx)^2 + (dy)^2 + (dz)^2 = - (cdt)^2 + (dr)^2 + (rd\theta)^2 + (r \sin \theta d\phi)^2$$

- Position of minus sign is just a choice.

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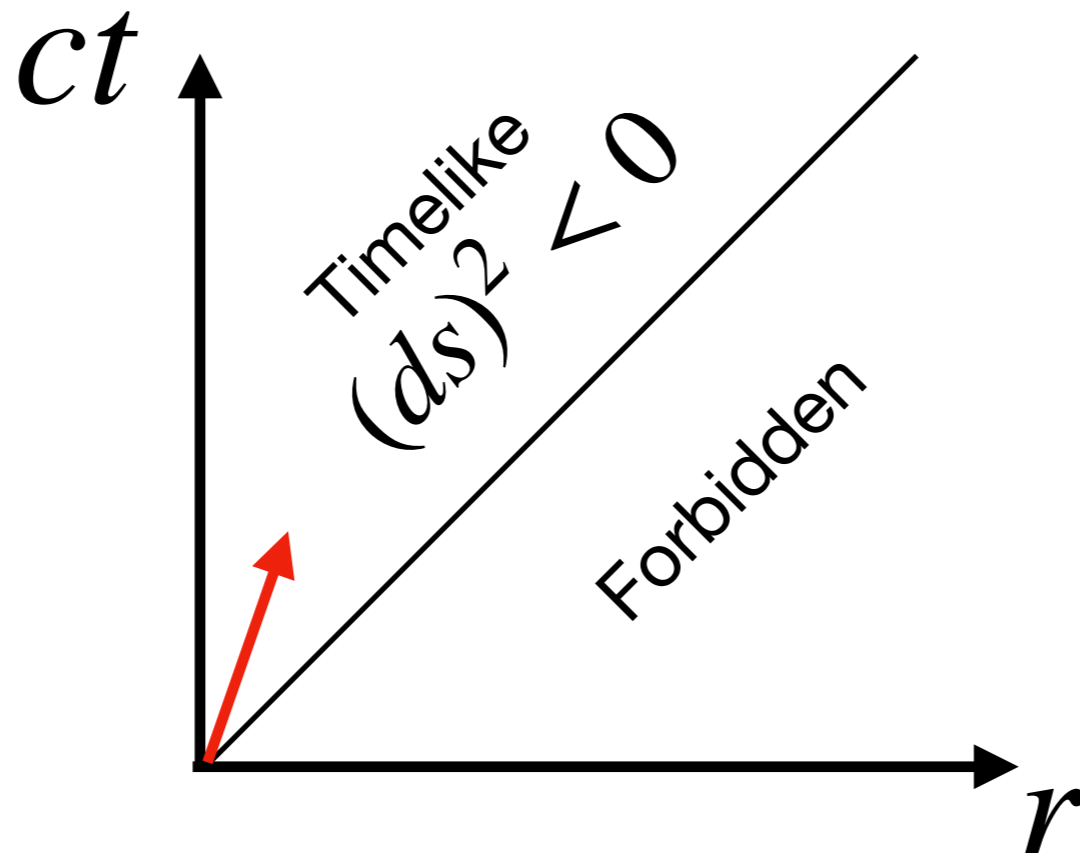
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- Can either have minus sign in front of coordinate time and proper time or in front of spatial coordinates.

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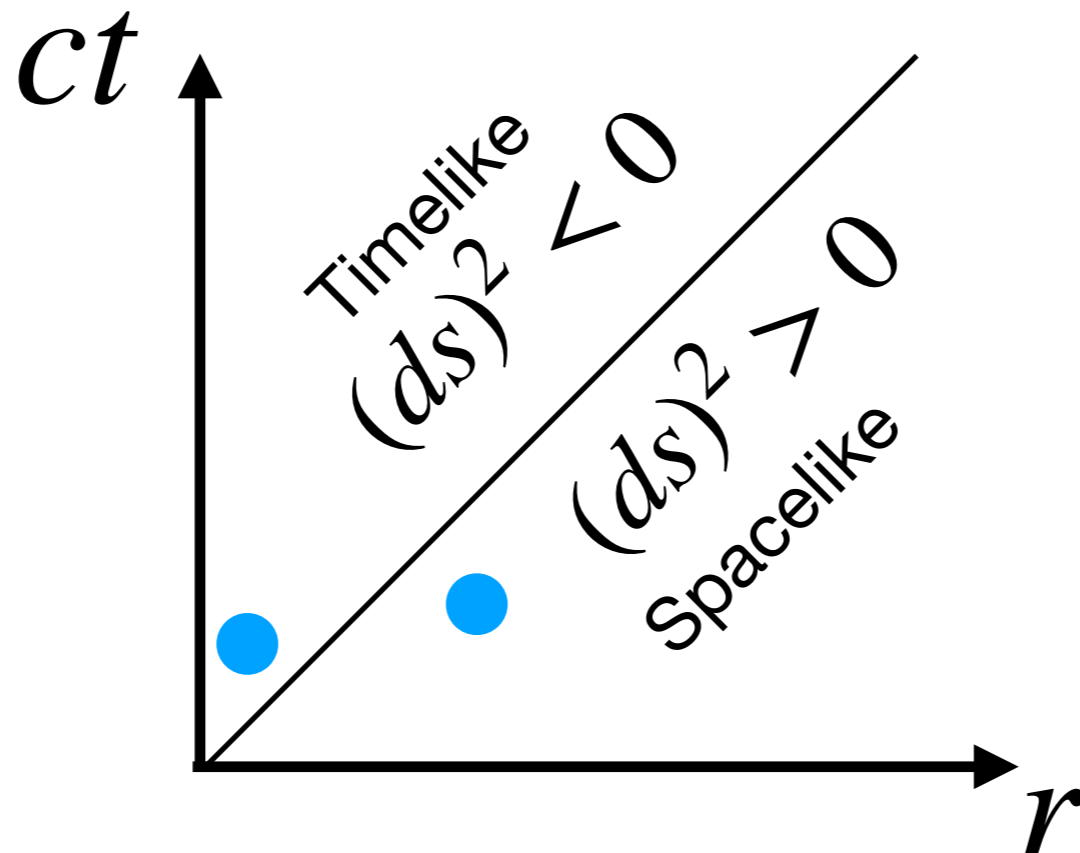
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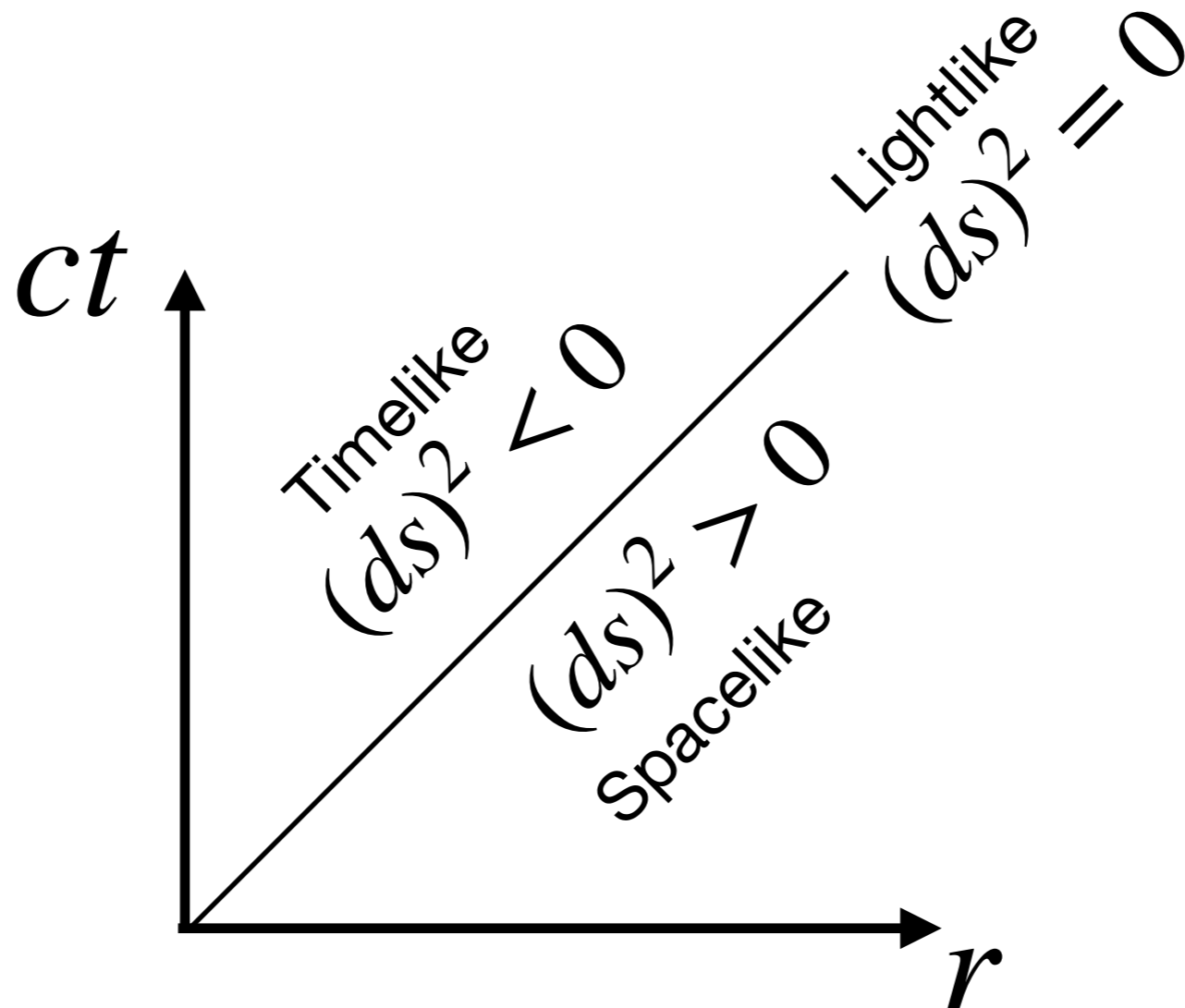
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- Light travels at c on null worldlines: $(ds)^2 = 0$



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- Cartesian:

$$\eta_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

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- Metric depends on spacetime curvature, but also on coordinate system. Line element (spacetime interval) is independent of coordinate system.

The Einstein Field Equations

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Einstein tensor

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Stress-energy tensor

= mass density and pressure
In SR, this is:

$$T_{\mu\nu} = \begin{pmatrix} \rho_0 c^2 & 0 & 0 & 0 \\ 0 & P_x & 0 & 0 \\ 0 & 0 & P_y & 0 \\ 0 & 0 & 0 & P_z \end{pmatrix}$$

The Schwarzschild Metric

$$G_{\mu\nu} = 8\pi \frac{G}{c^4} T_{\mu\nu}$$

- Solve Einstein equations for spherically symmetric spacetime with $\rho(\mathbf{r}) = M \delta^3(\mathbf{r})$
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- Tends to Minkowski for large r , but for small r angles of a triangle no longer add up to 180 degrees!

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- The same as our naive Newtonian guess!

The Innermost Stable Circular Orbit

- Now let's do orbits in the Schwarzschild metric.
- Energy equation:

$$\text{KE per unit mass} + \text{PE per unit mass} = \text{Total energy per unit mass}$$

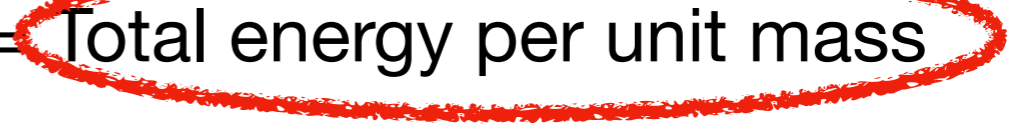
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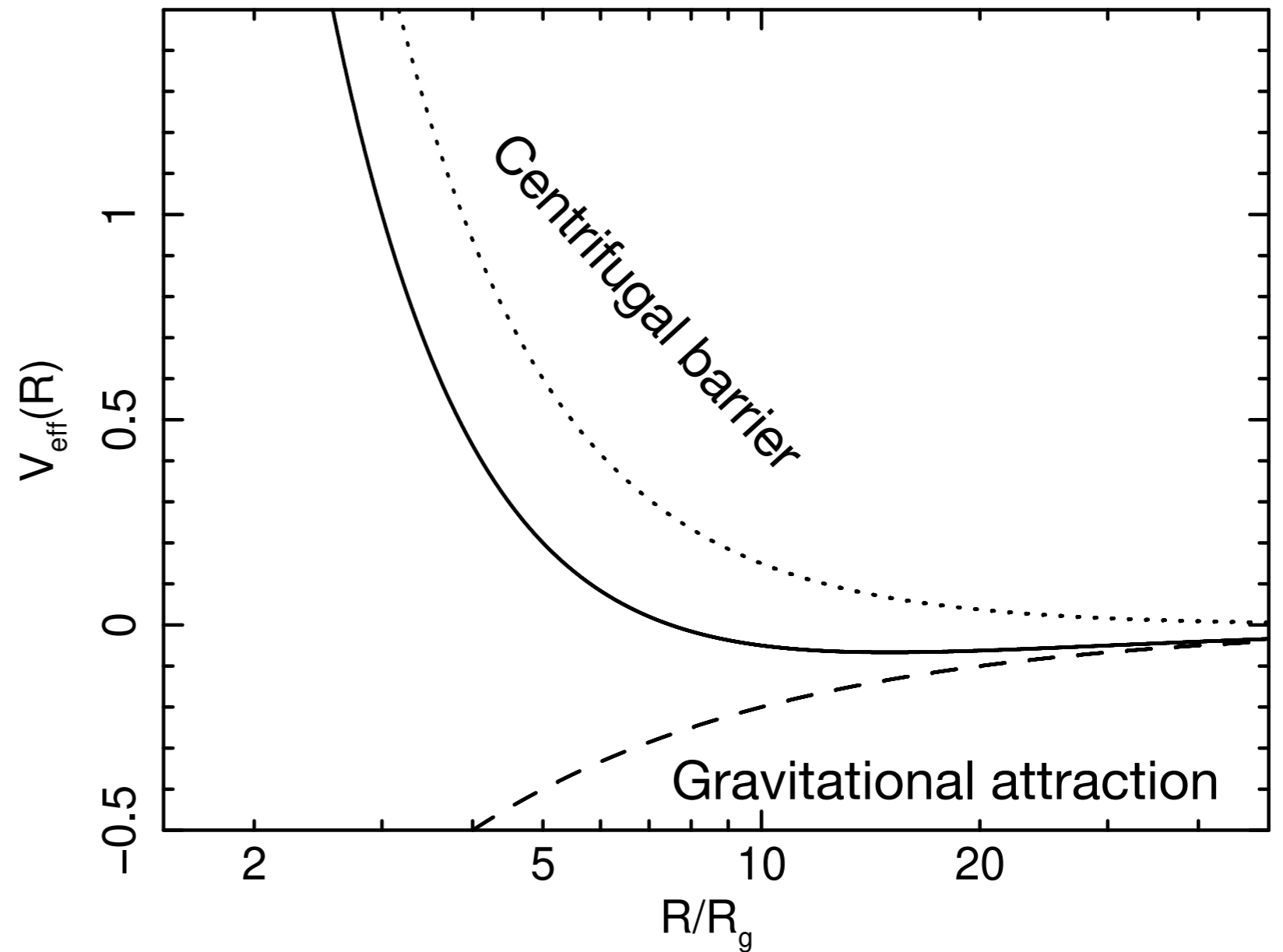
Centrifugal barrier

Gravitational attraction

- Circular orbit: $dr/dt=0$, so $V(r)=E$.
- $E=\text{constant}$, so for a circular orbit $dV(r)/dr=0$: circular orbits at turning points of $V(r)$.
- Minima = stable, maxima and inflection points = unstable.

The Innermost Stable Circular Orbit

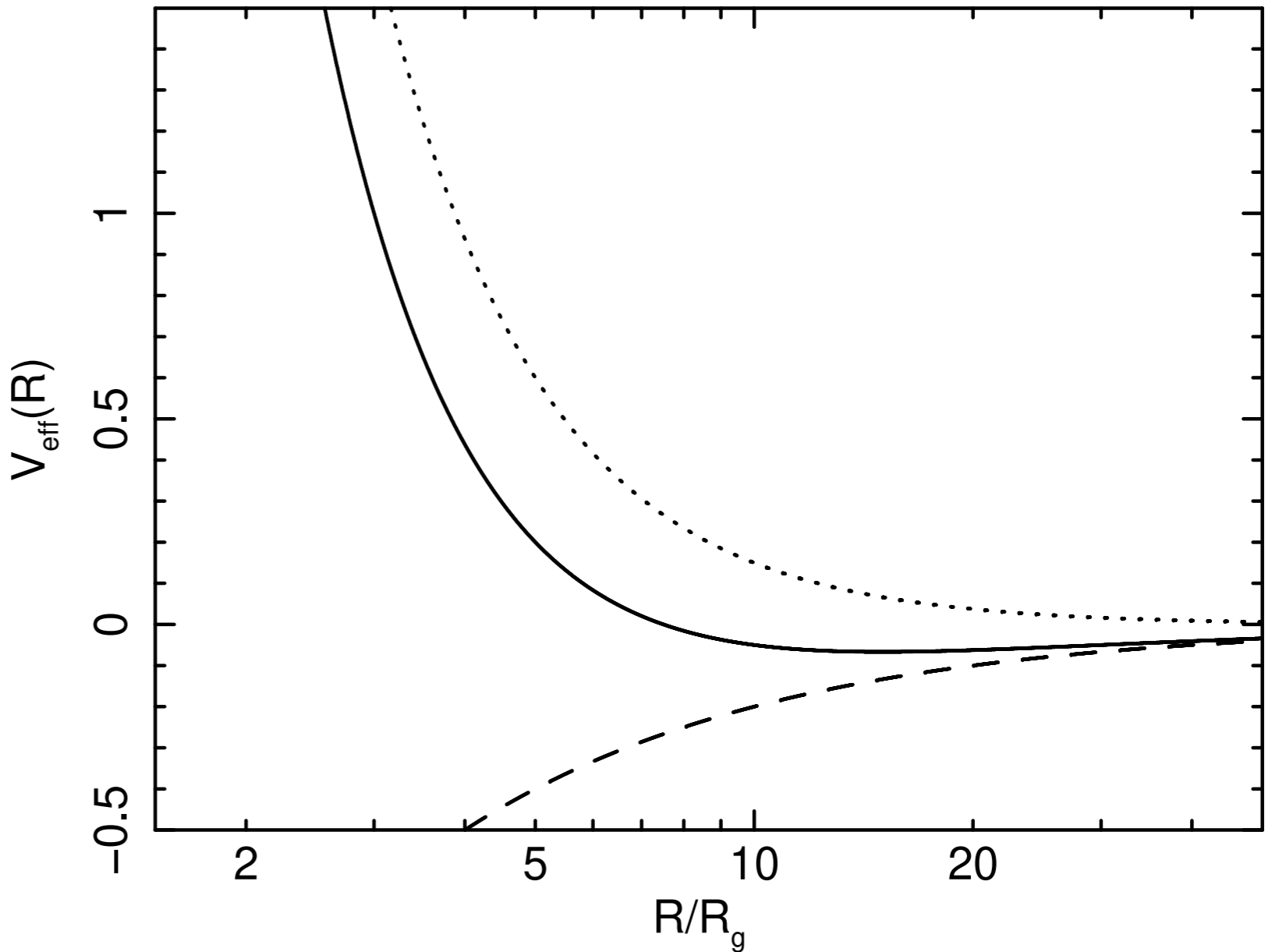
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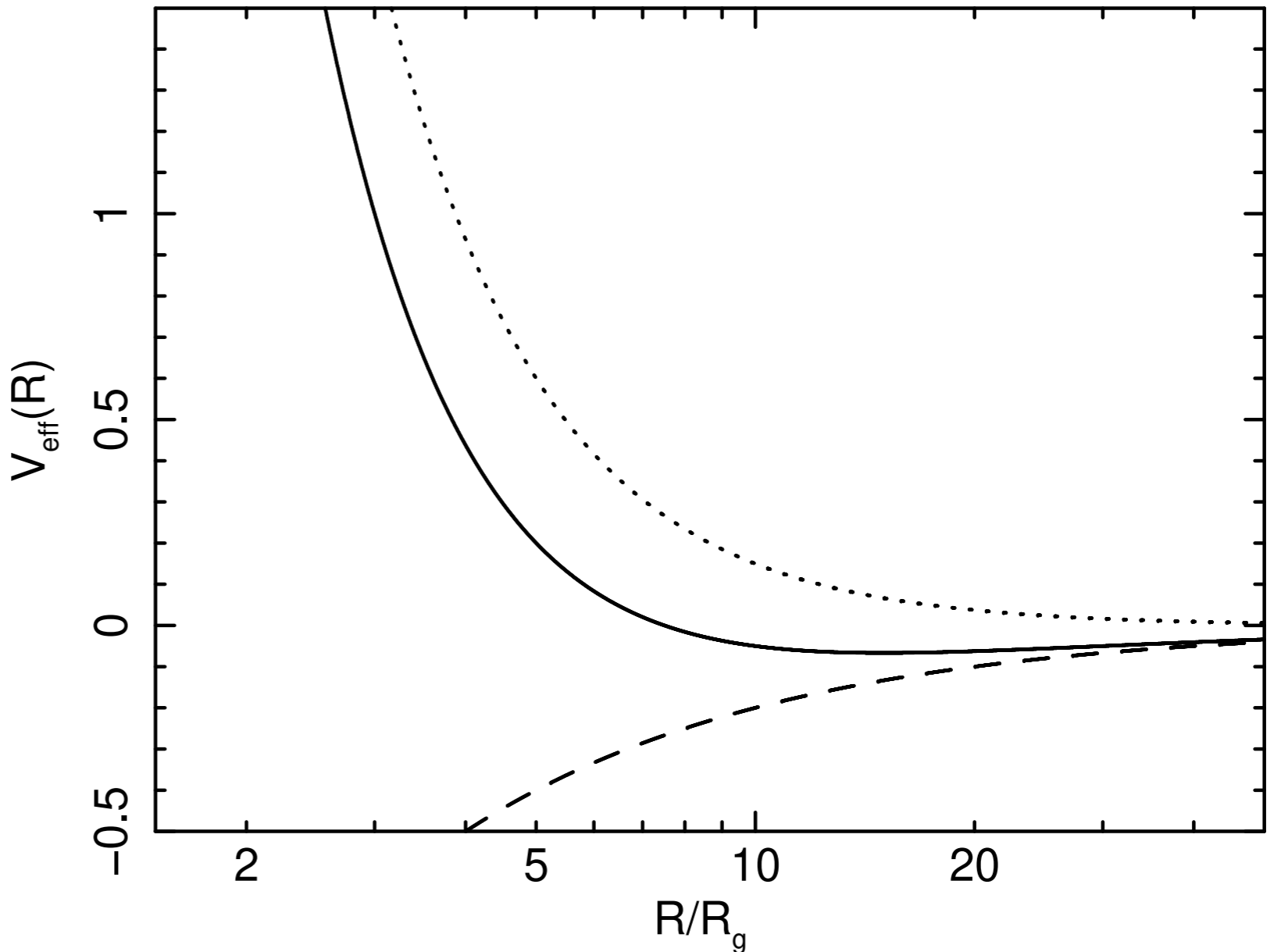


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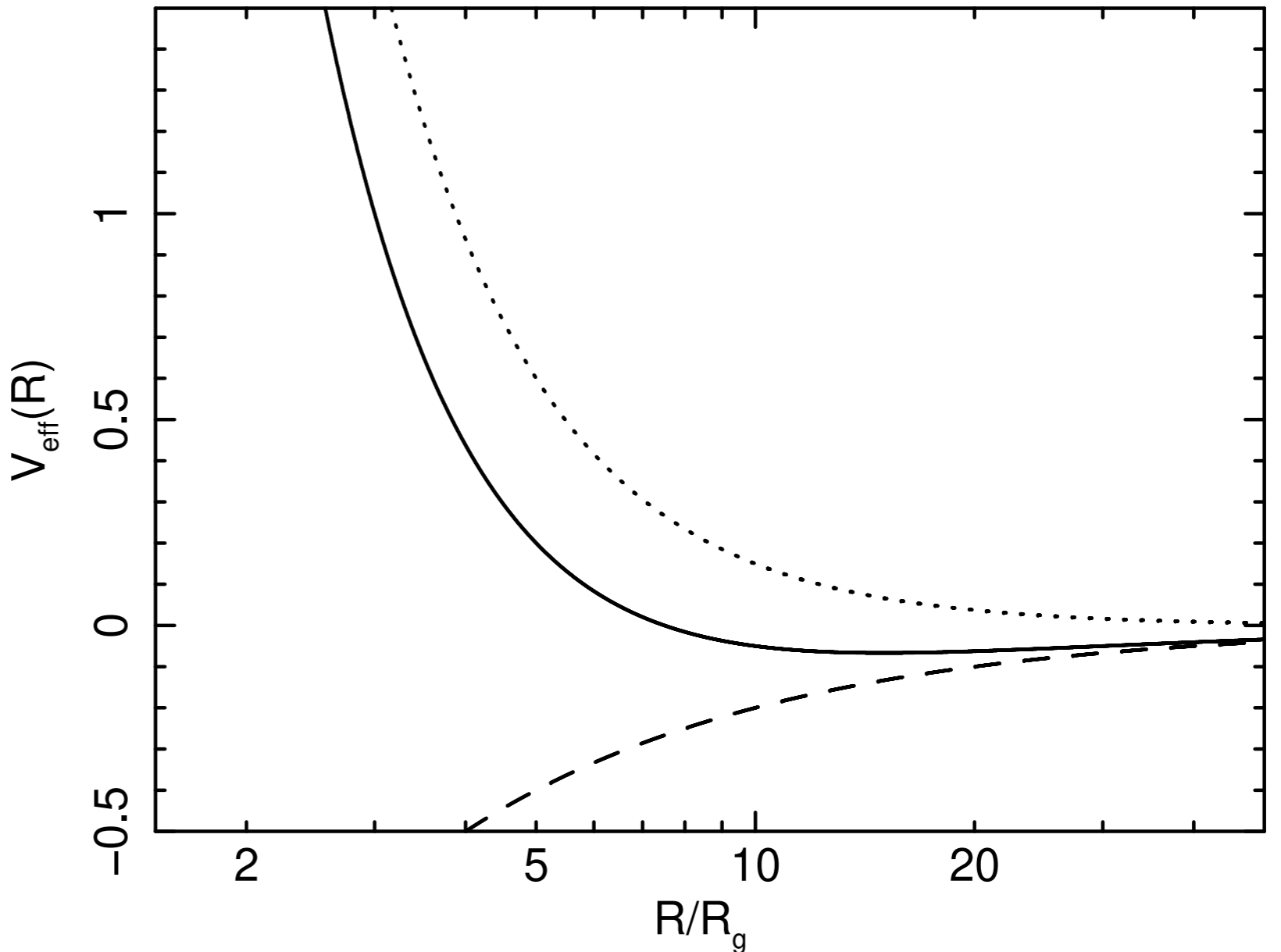


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$$L = vr \quad \therefore v^2 r^2 = GMr \quad \therefore v^2 = GM/r$$

...Keplerian orbit!

The Innermost Stable Circular Orbit

- For Schwarzschild solution, energy equation becomes:

$$\left(\frac{dr}{d\tau}\right)^2 + V_{\text{eff}}^2(r) = \left(\frac{E}{c}\right)^2$$

- With effective potential:

$$V_{\text{eff}}^2(r) = \left(1 - \frac{2r_g}{r}\right) \left(\frac{L^2}{r^2} + c^2\right)$$

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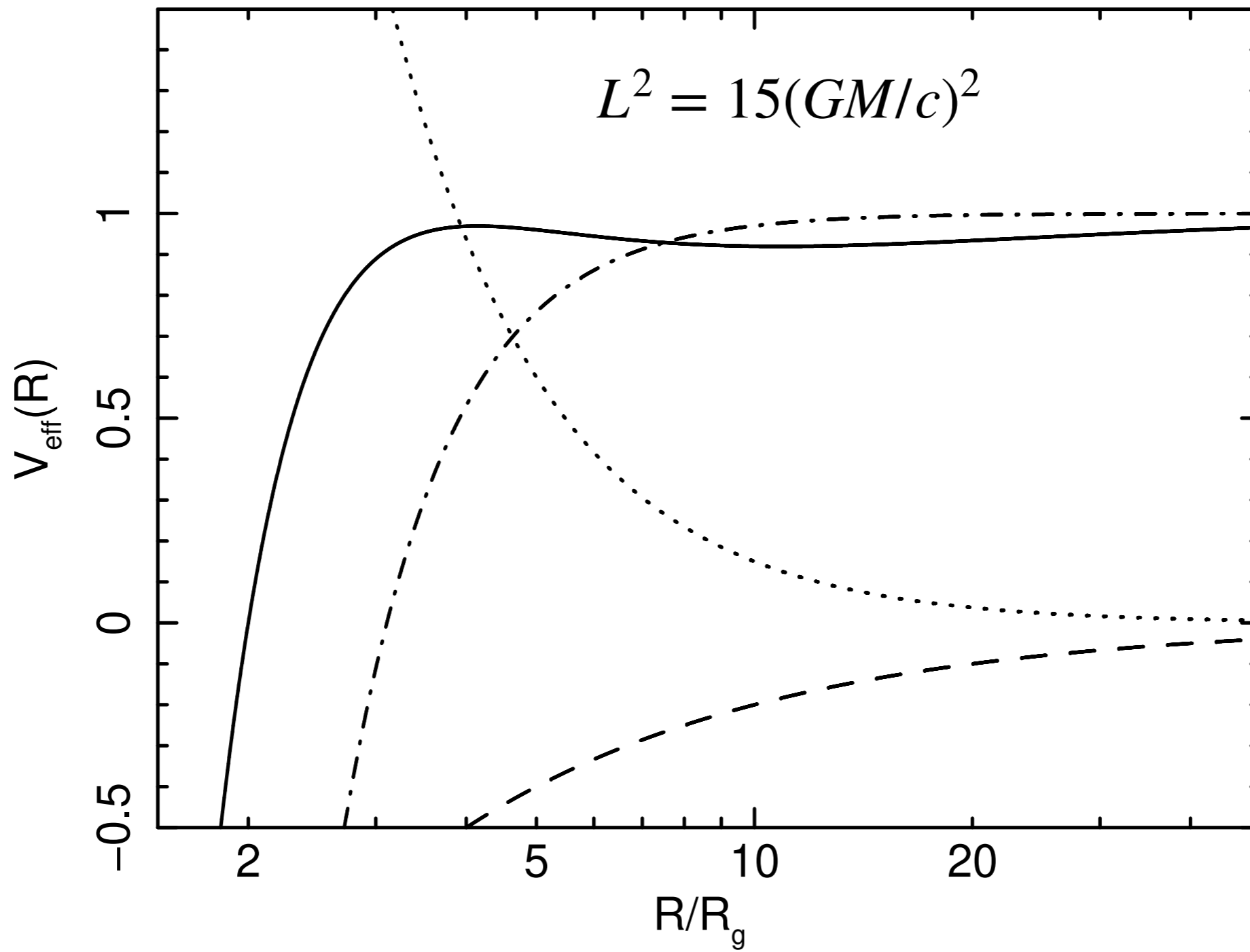
Centrifugal barrier

Gravitational attraction

Completely new!

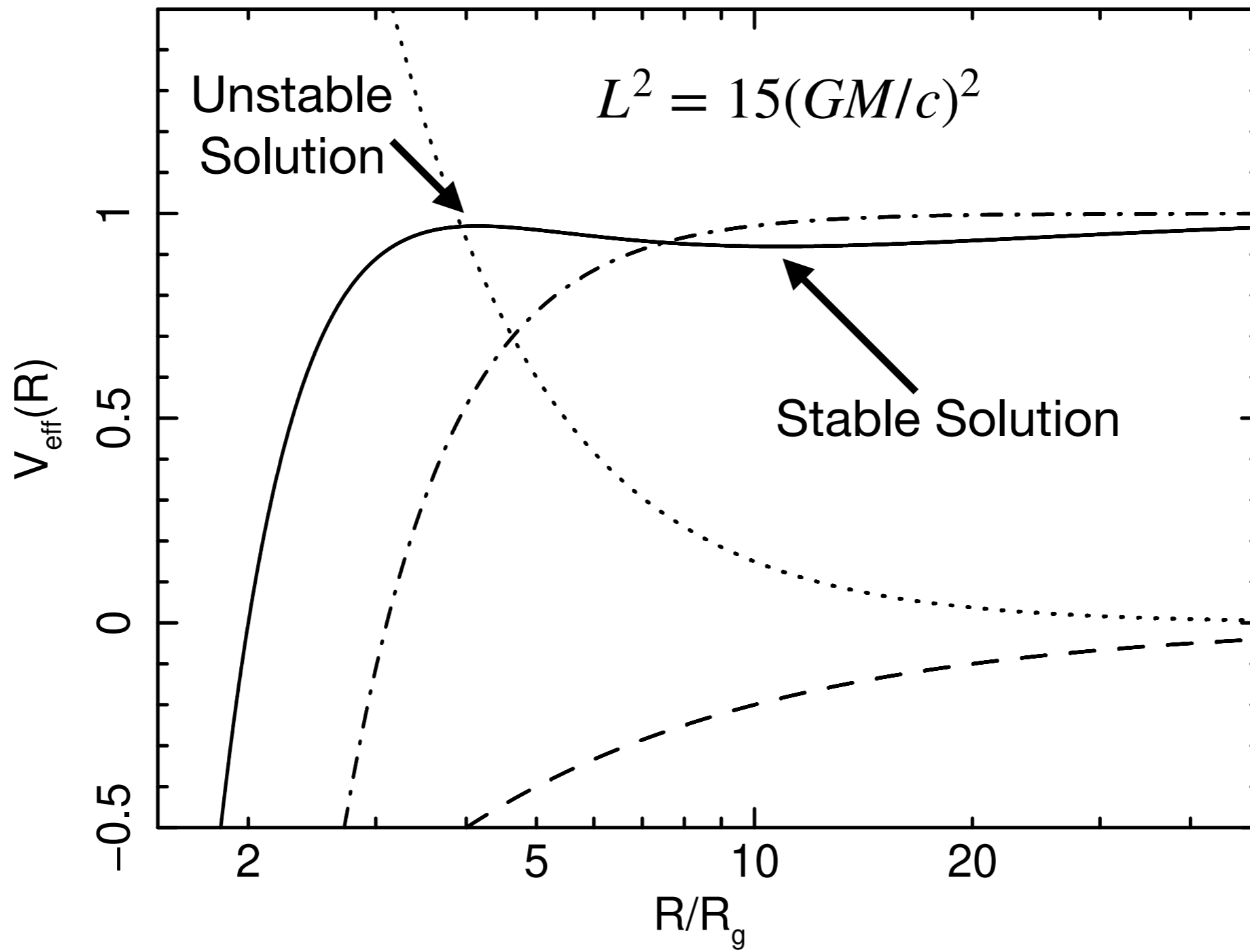
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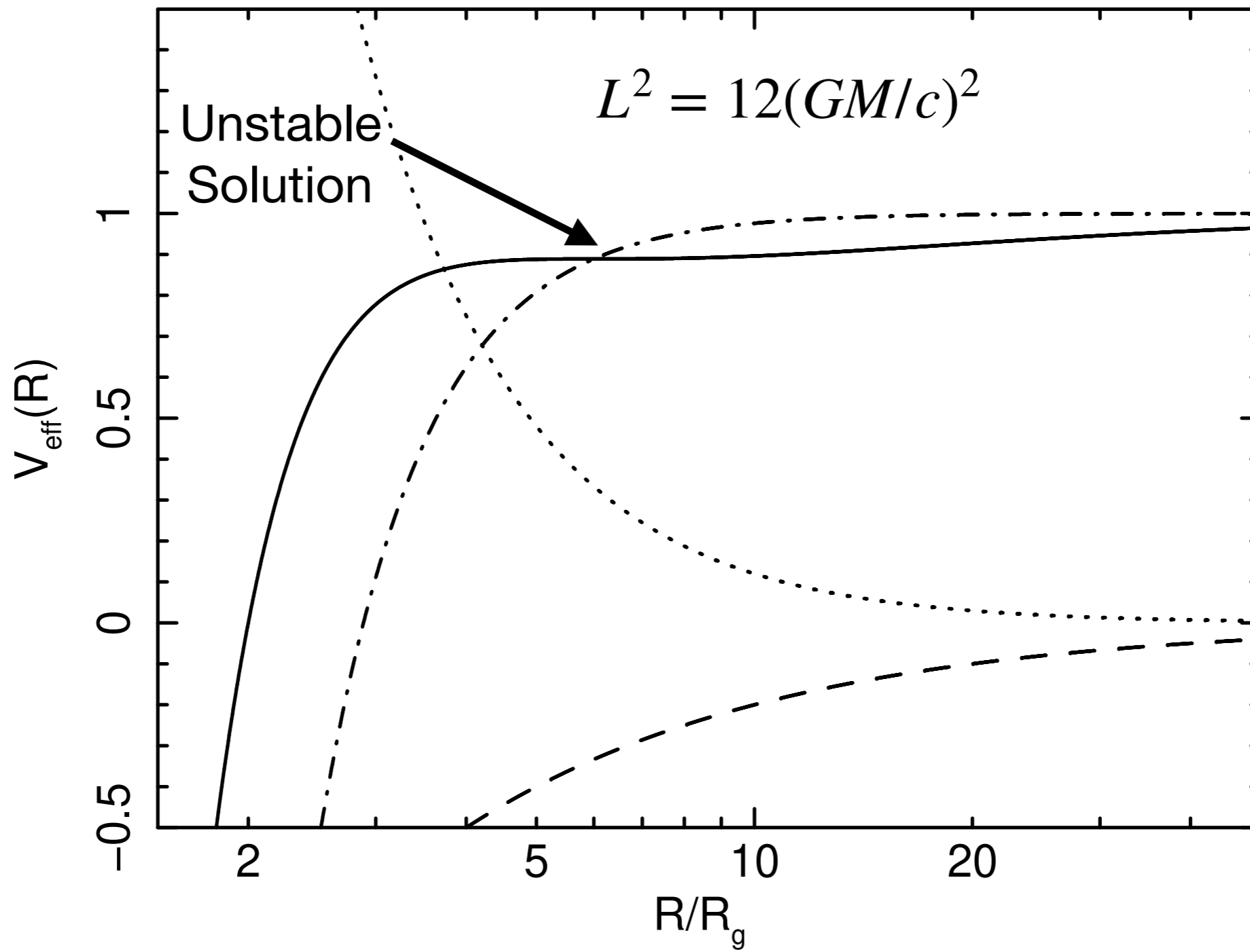
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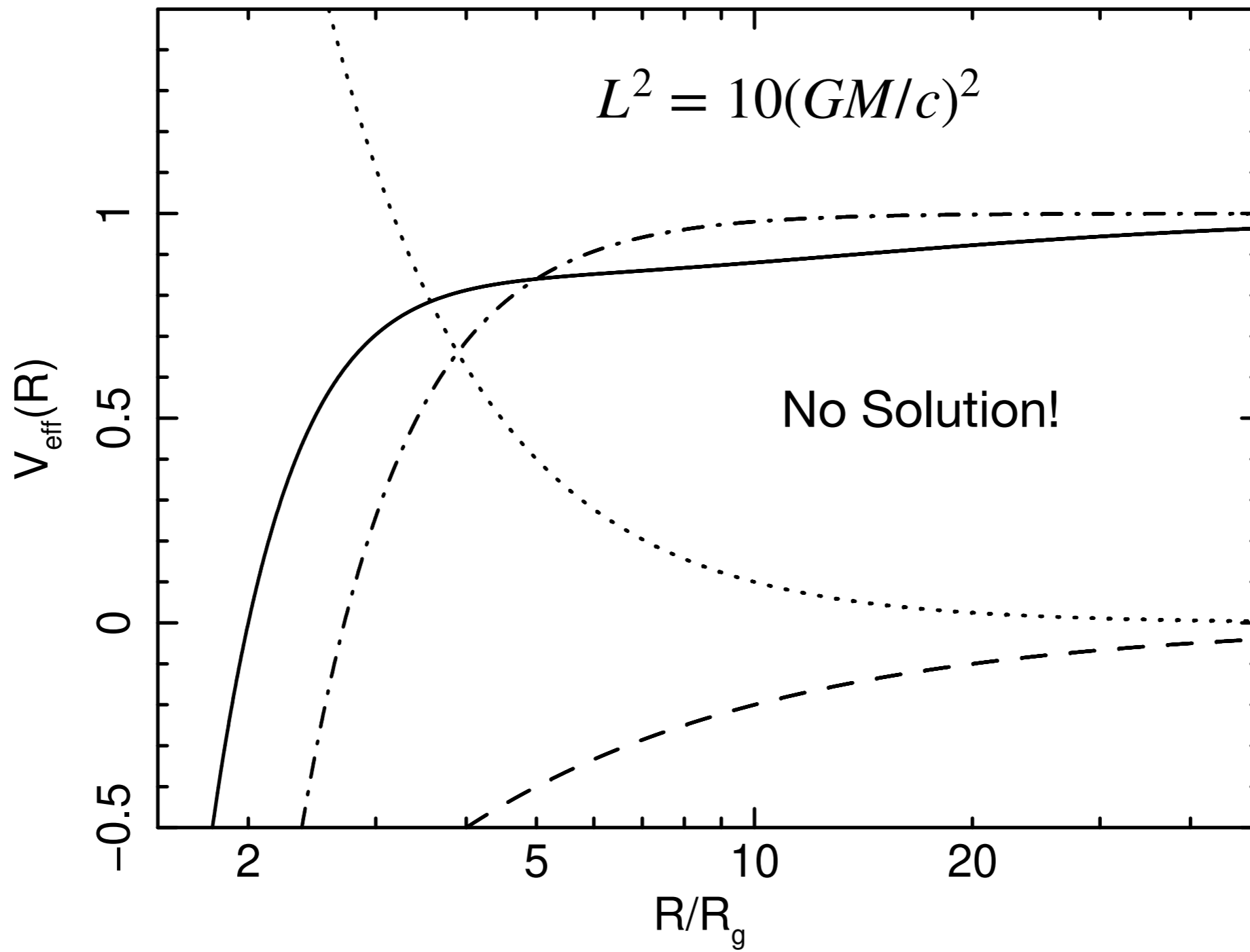
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$$r = \frac{L^2}{2GM} \left[1 \pm \sqrt{1 - 12 \left(\frac{GM}{cL} \right)^2} \right]$$

...two solutions:
stable and unstable!

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- No solutions for: $L < L_{\text{crit}}$
- Therefore ISCO at:

$$r_{\text{isco}} = \frac{L_{\text{crit}}^2}{2GM} = \frac{12(GM)^2}{2GMc^2} = 6 r_g$$

- Inside $6 r_g$, can still in principle escape the BH, but can't orbit without help!

The Kerr Metric

- If the BH is spinning, no longer have spherical symmetry, only azimuthal.
- Kerr solution (Kerr 1960) is:

$$g_{tt} = - \left(1 - \frac{2r}{r_g \Sigma} \right) \quad g_{t\phi} = g_{\phi t} = \frac{2ar \sin^2 \theta}{r_g \Sigma} \quad g_{rr} = \frac{\Sigma}{\Delta} \quad g_{\theta\theta} = \Sigma \quad g_{\phi\phi} = \frac{\mathcal{A} \sin^2 \theta}{\Sigma}$$

$$\Sigma = (r/r_g)^2 + a^2 \cos^2 \theta \quad \Delta = (r/r_g)^2 - 2(r/r_g) + a^2 \quad \mathcal{A} = [(r/r_g)^2 + a^2]^2 - \Delta a^2 \sin^2 \theta$$

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- a is the dimensionless spin parameter: $a = \frac{J_{\text{bh}}}{Mcr_g}$
- Boyer-Lindquist coordinates:

$$x = r_g \sqrt{(r/r_g)^2 + a^2} \sin \theta \cos \phi$$

$$y = r_g \sqrt{(r/r_g)^2 + a^2} \sin \theta \sin \phi$$

$$z = r \cos \theta$$

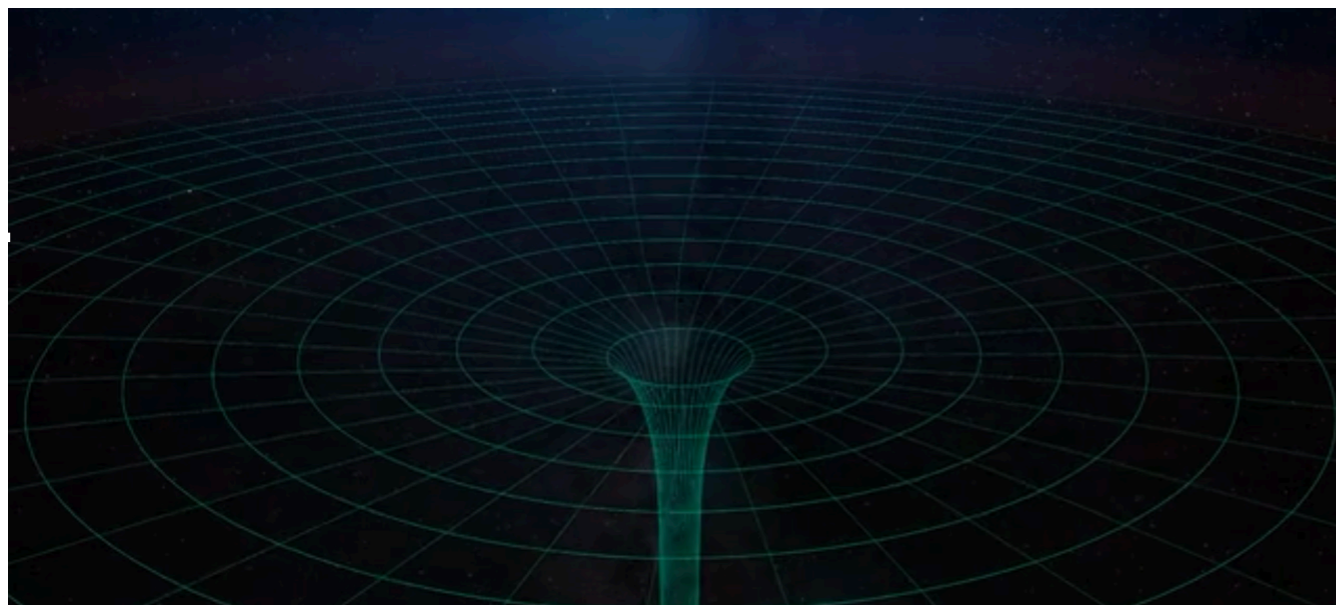
The Kerr Metric

- Cross term = **Frame Dragging Effect**

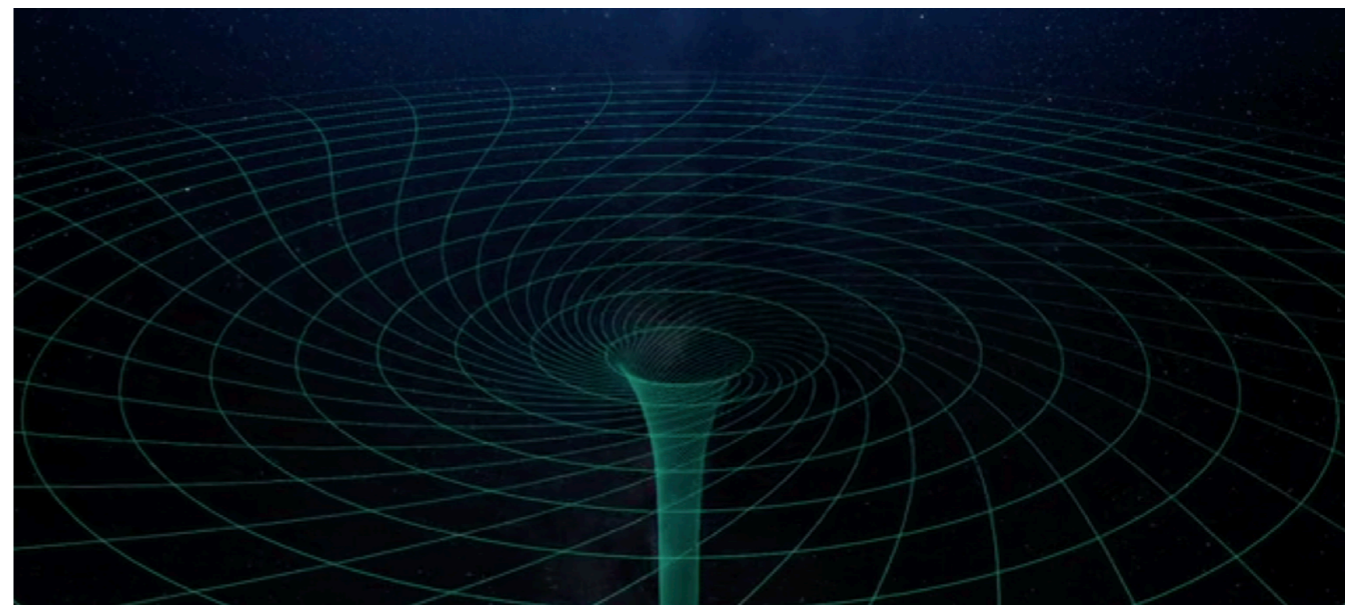
$$g_{t\phi} = g_{\phi t} = \frac{2ar \sin^2 \theta}{r_g \Sigma}$$

- Kerr BH drags spacetime around with it.
- Gives rise to (among other things) Lense-Thirring precession — a vertical wobble of orbits in a plane inclined to the BH equatorial plane.

a = 0



a > 0



The Kerr Metric

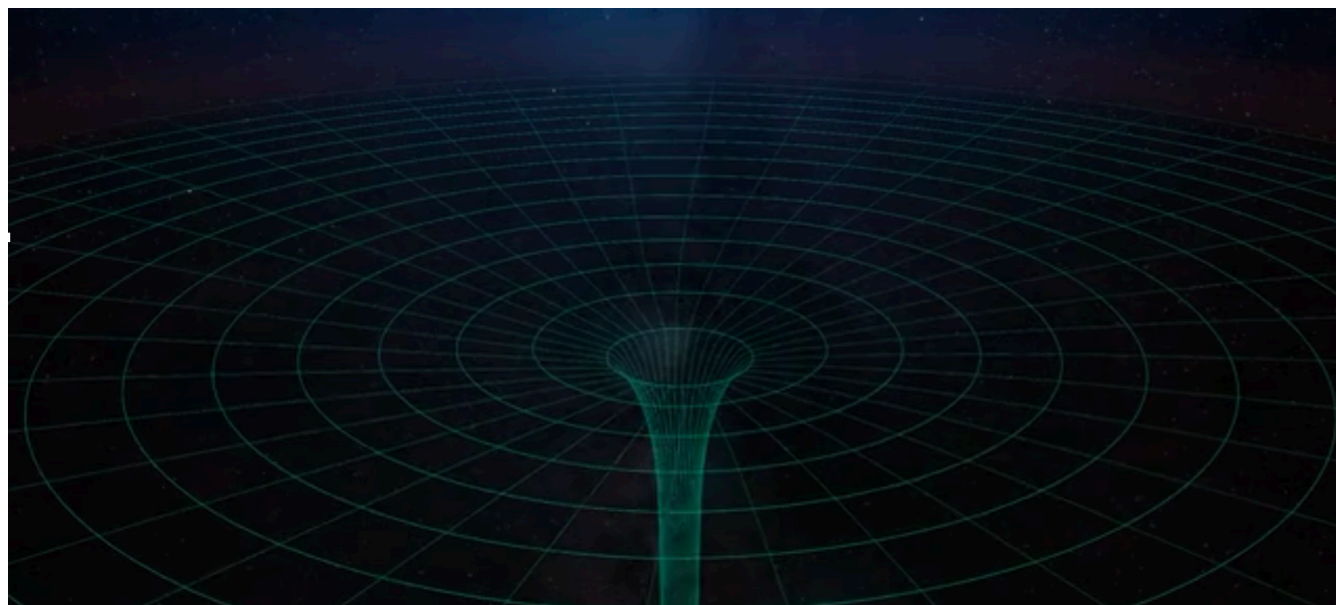
- Horizon is still at coordinate singularity of g_{rr} :

$$g_{rr} = \frac{\Sigma}{\Delta}$$

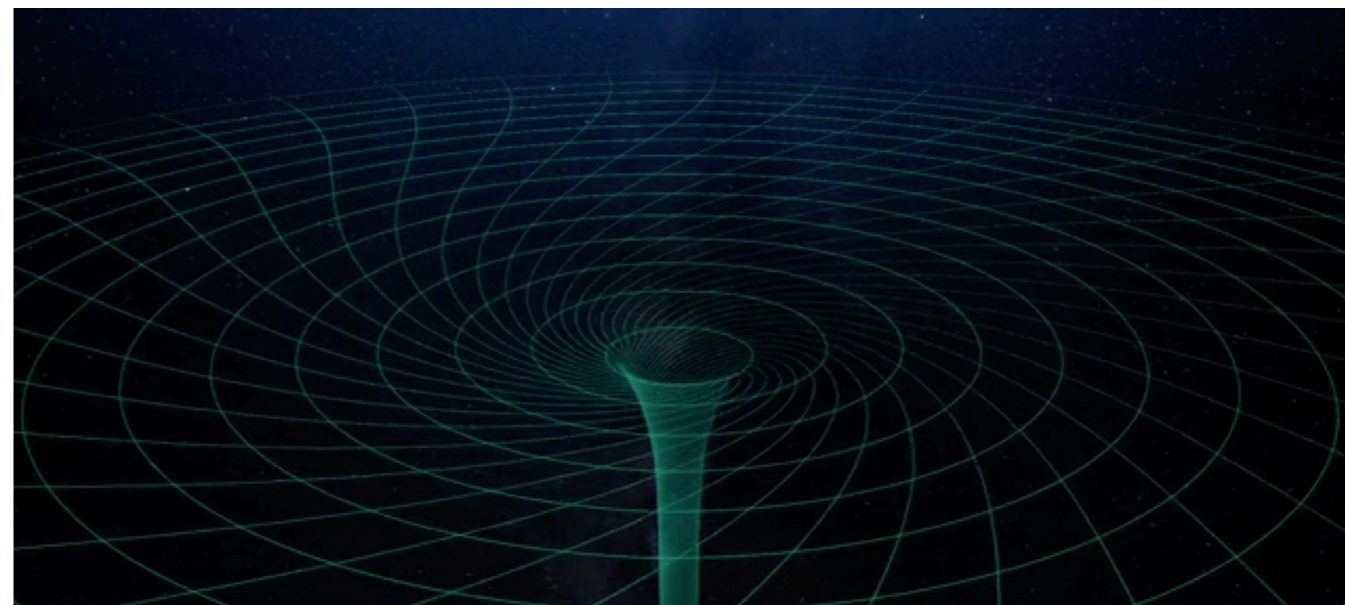
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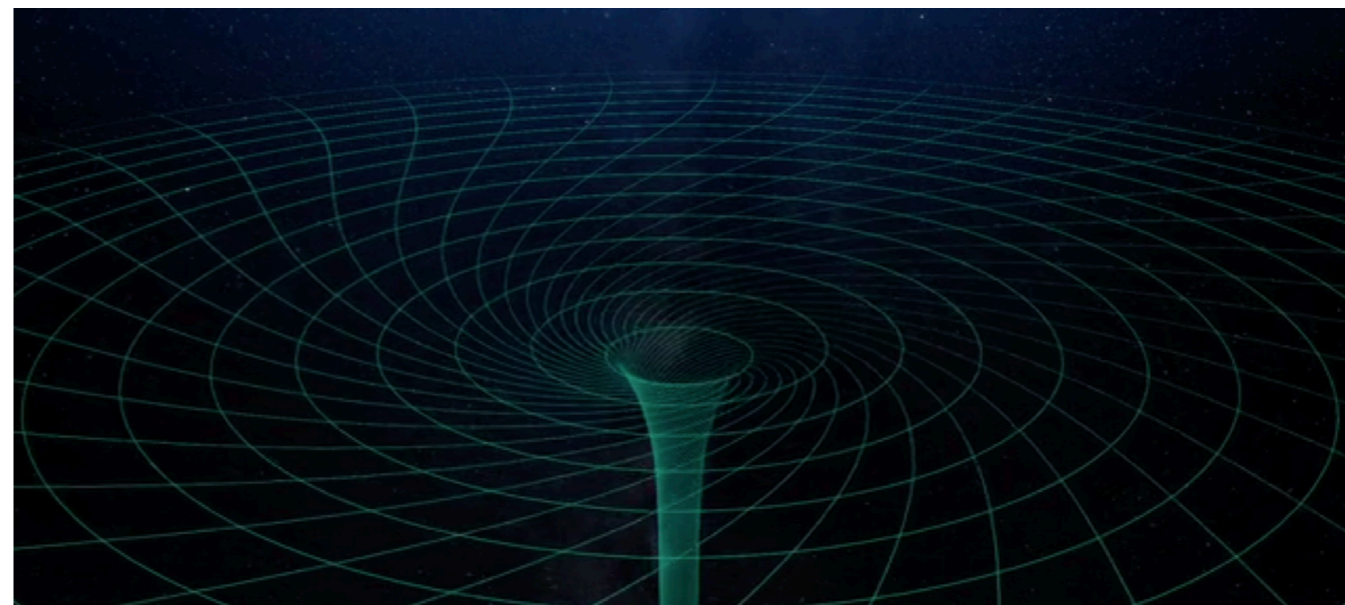
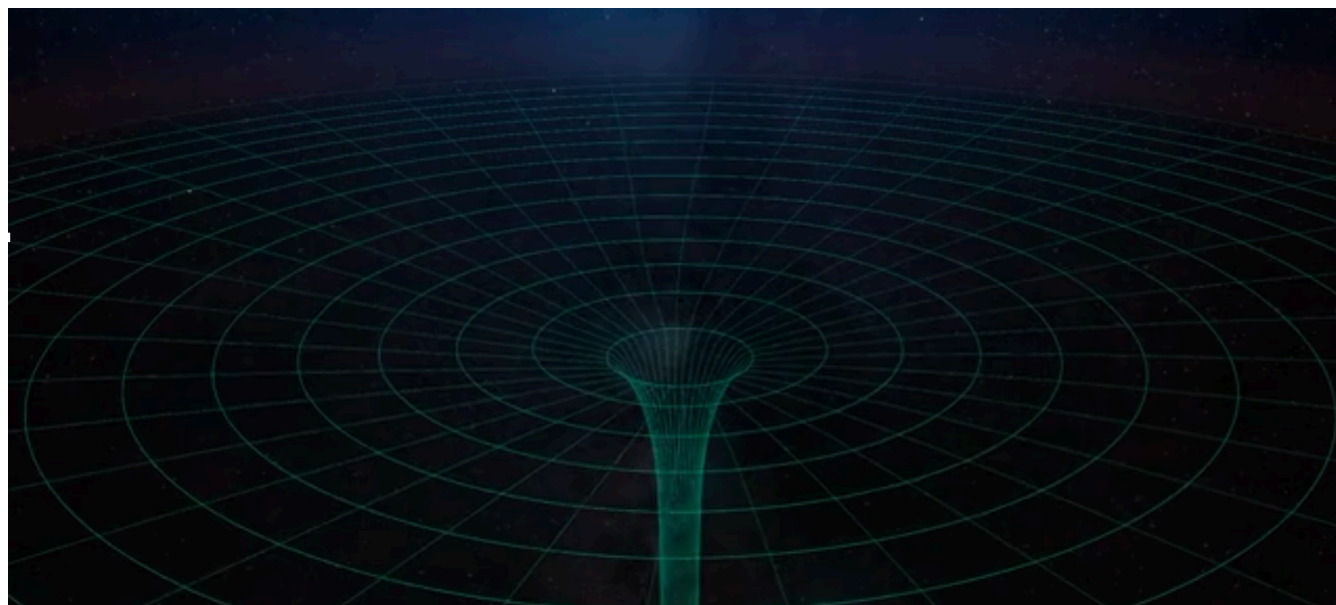
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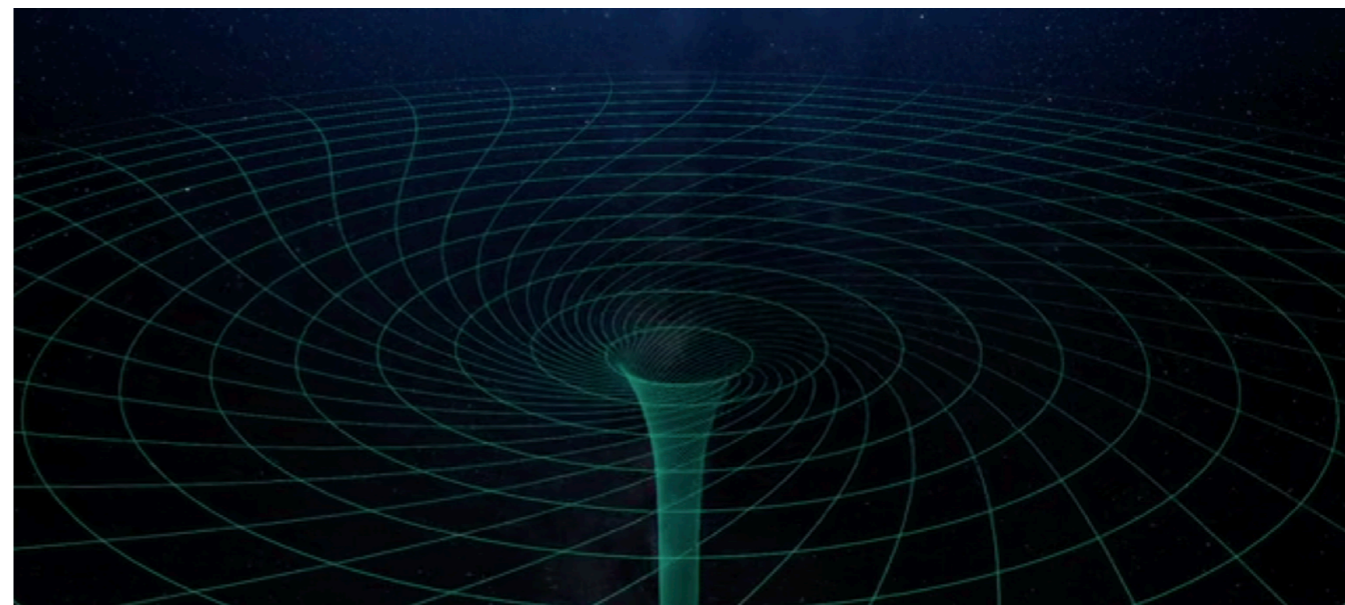
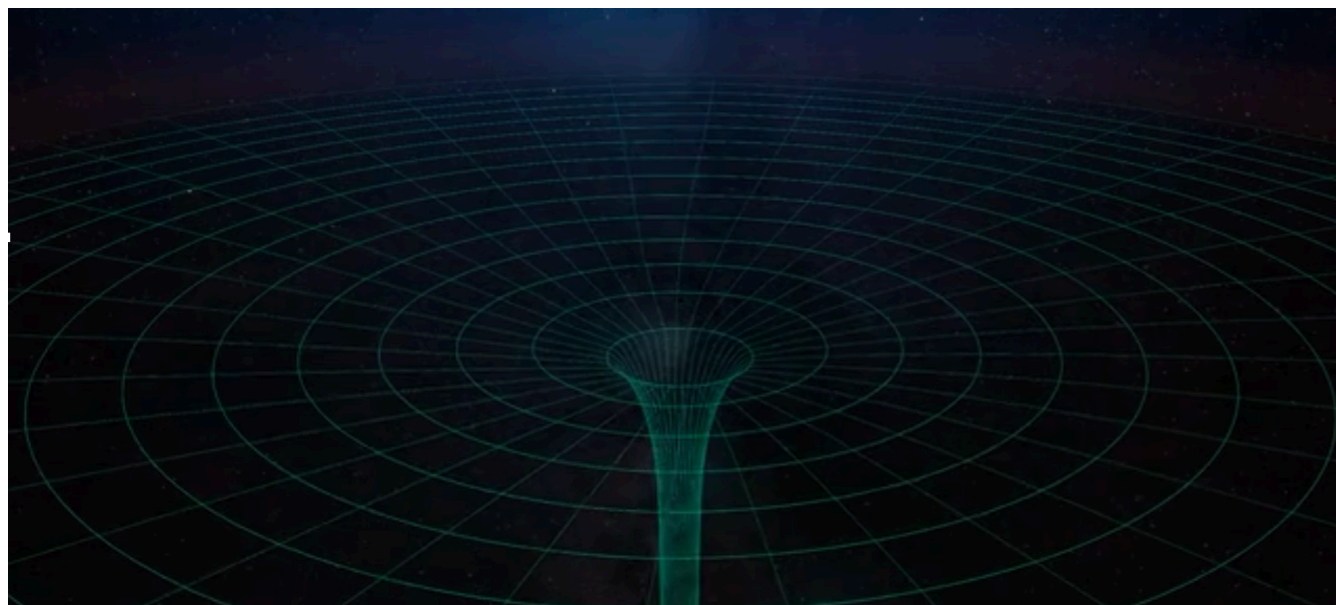
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- Causality: $-1 \leq a \leq 1$ where +ve is prograde and -ve is retrograde.

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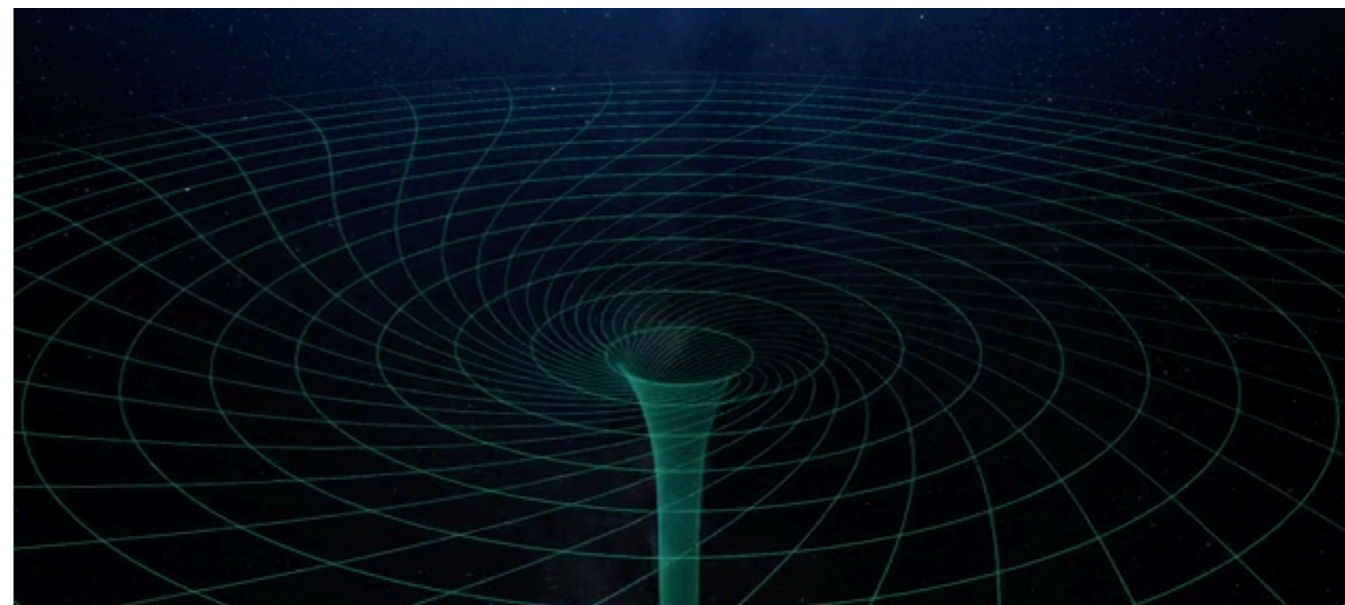
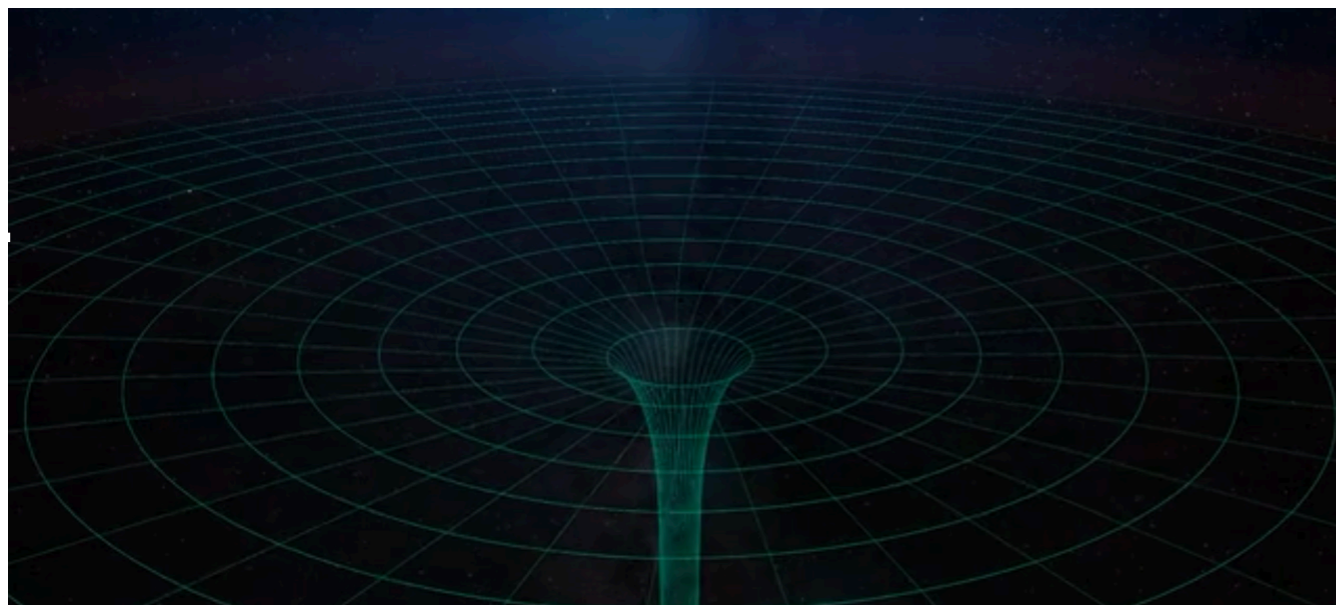
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- Can think of this like a rotating floor giving angular momentum in one direction and taking in the other.

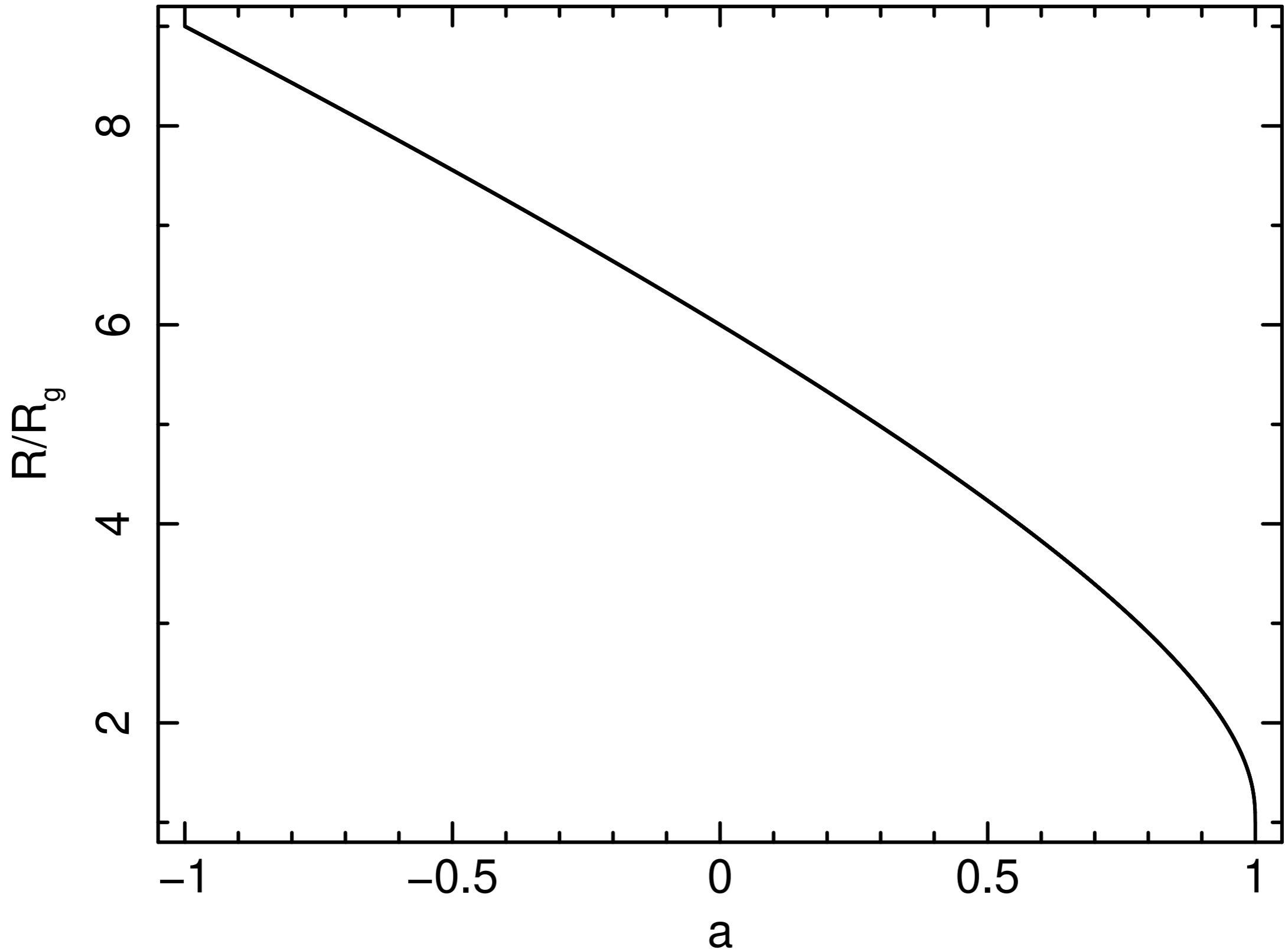
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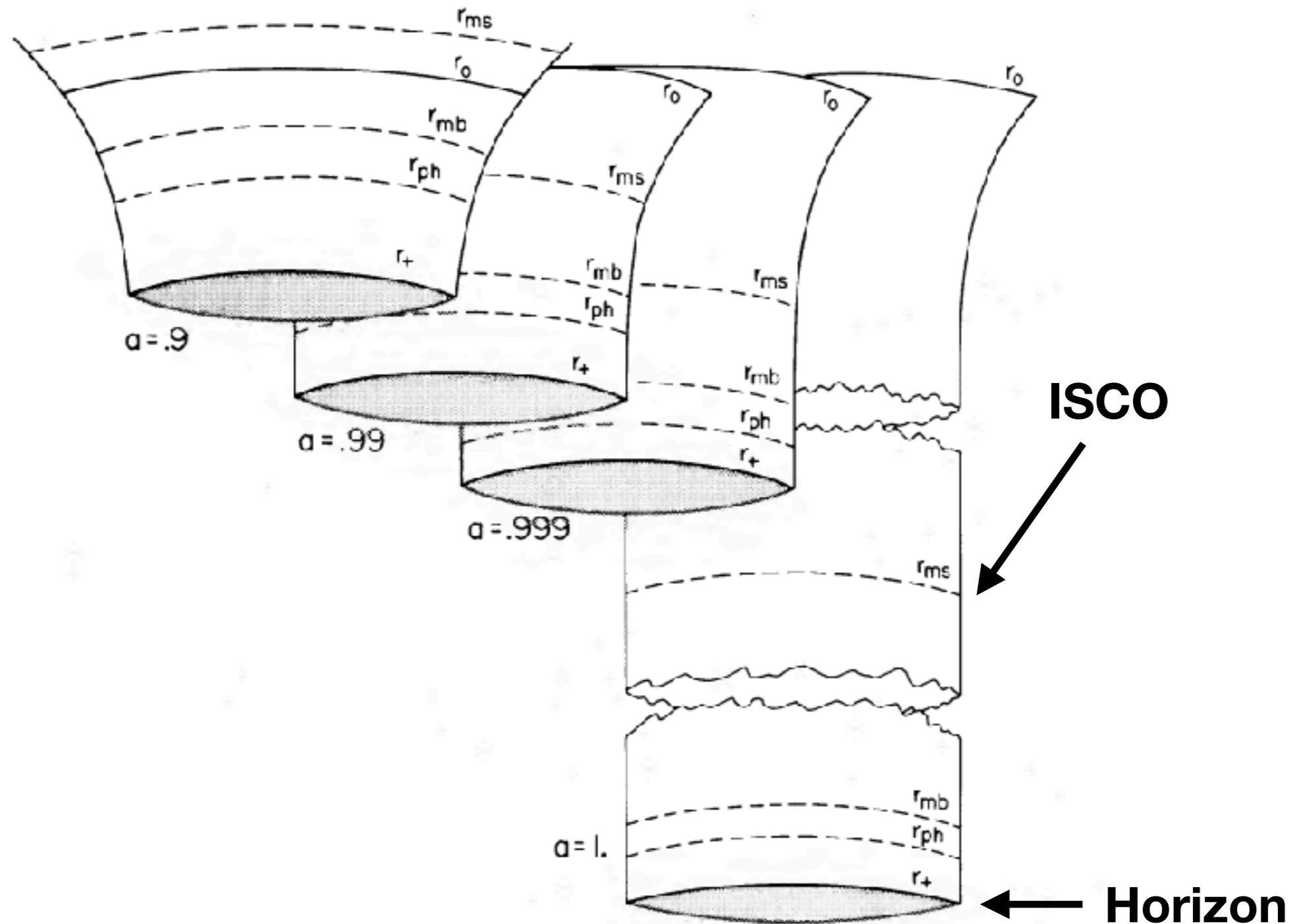
The Kerr Metric

- ISCO changes with spin:



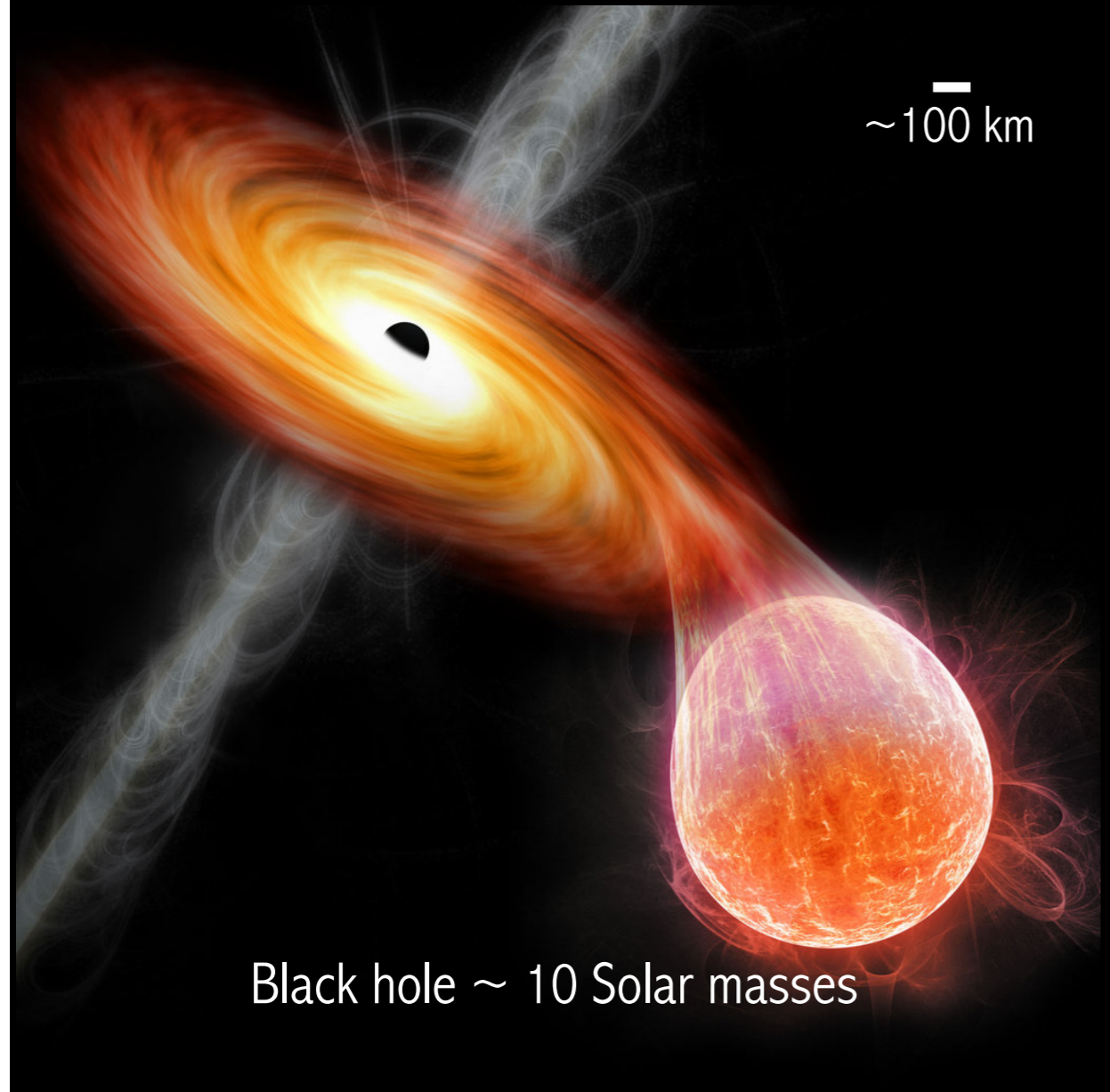
The Kerr Metric

- What's going on at $a=1$? ISCO and horizon both at $1r_g$, but not in the same place! Break down of B-L coordinates.

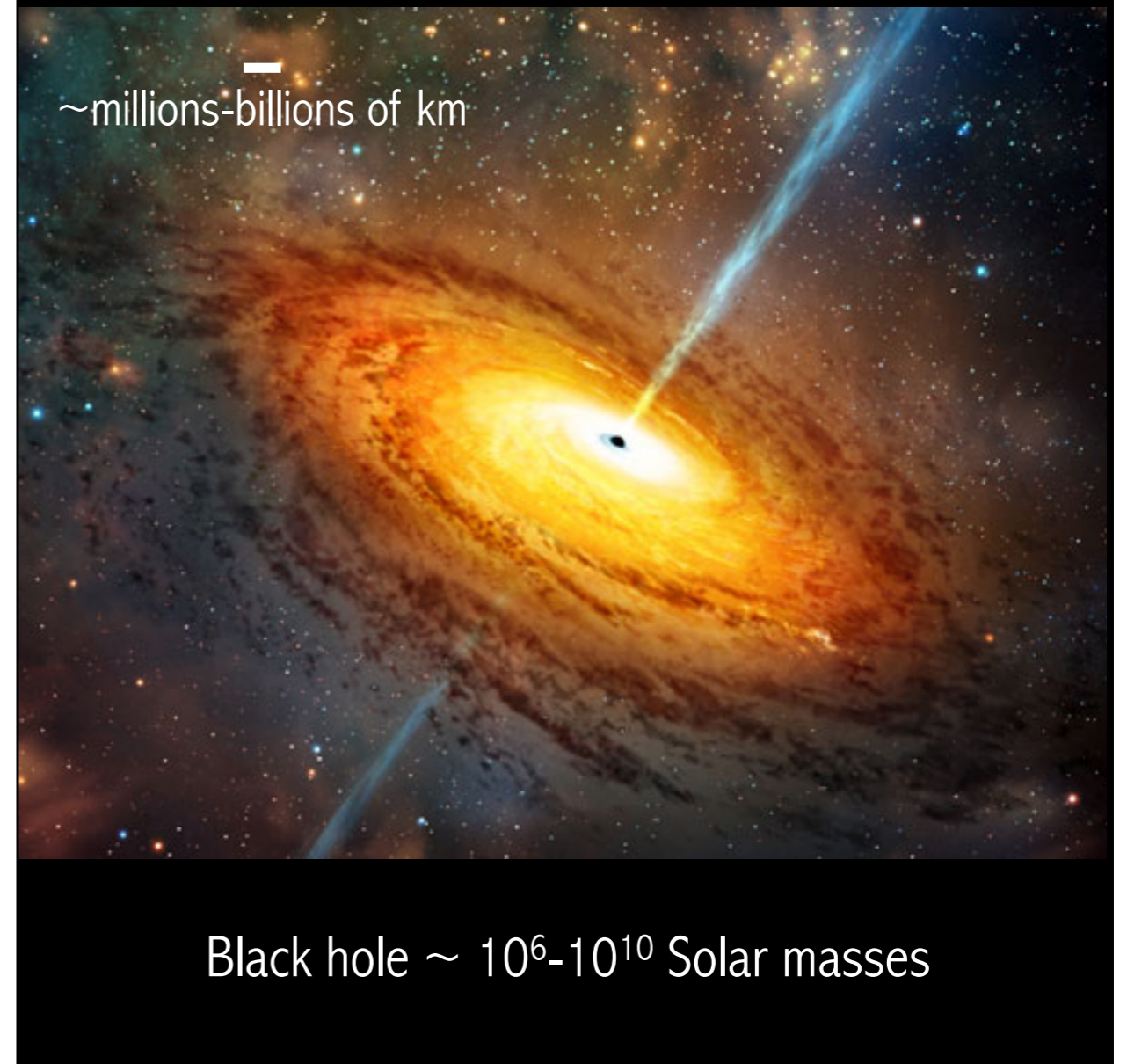


Accreting Black Holes

Black Hole X-ray Binaries

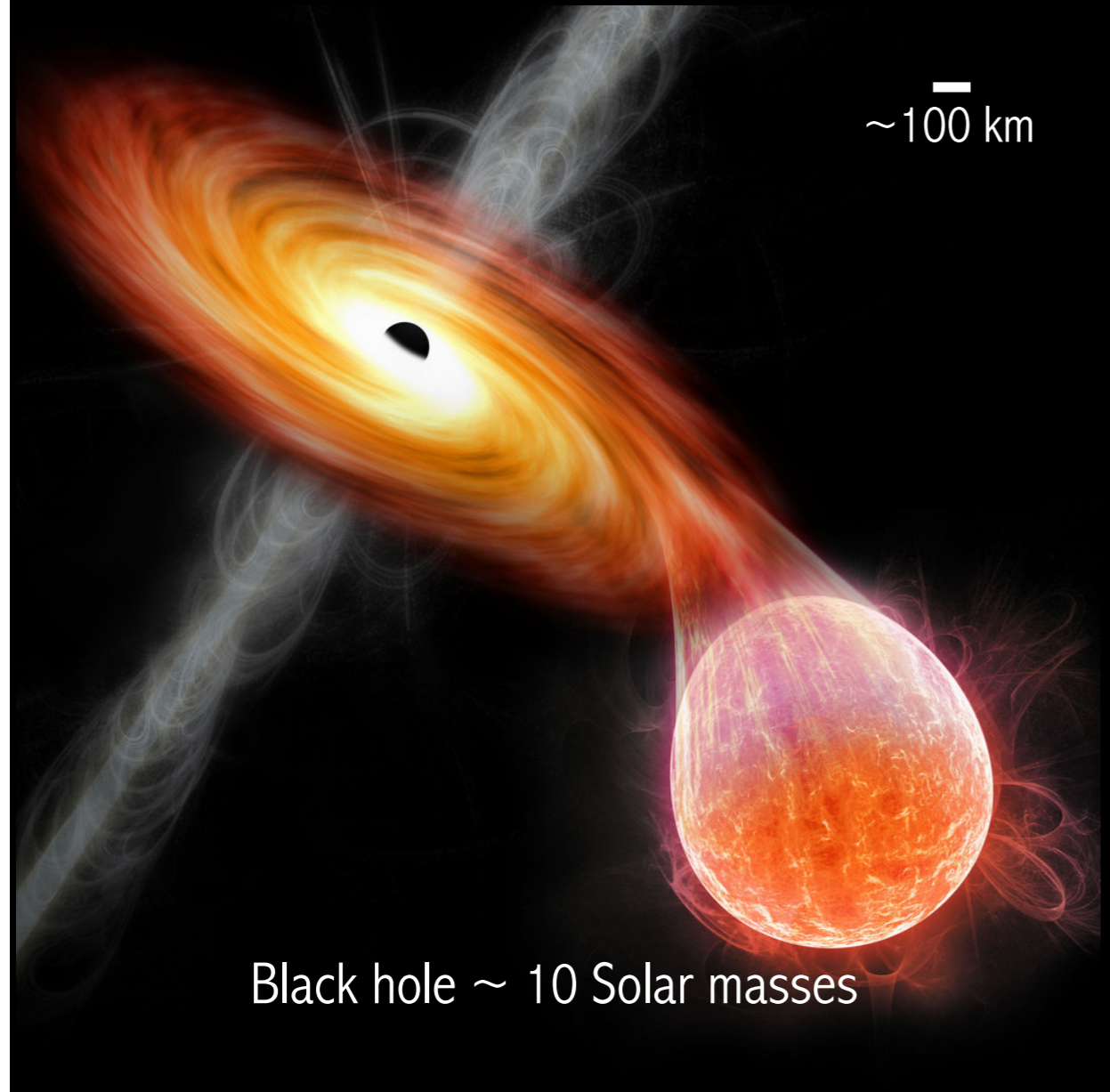


Active Galactic Nuclei

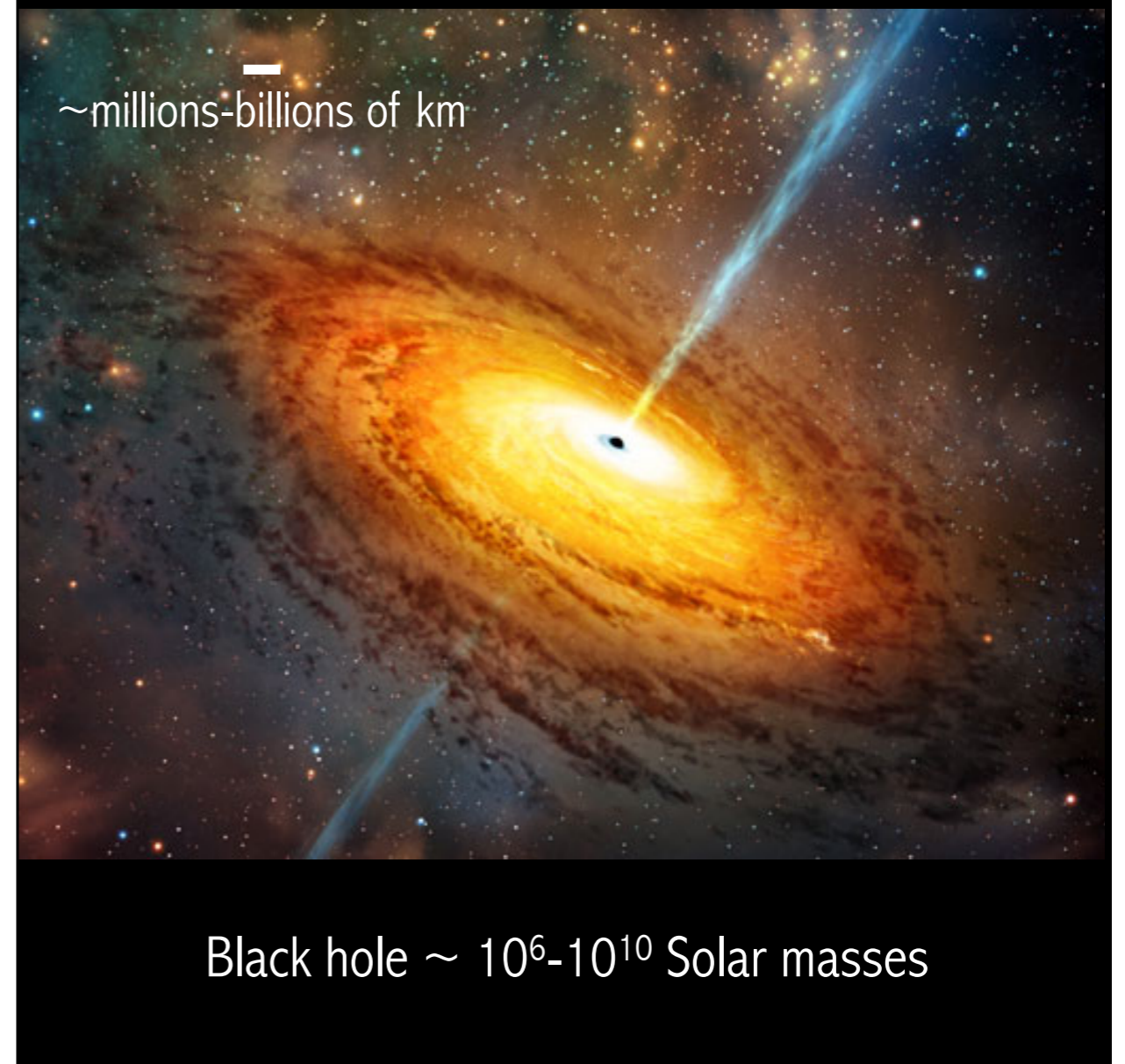


Accreting Black Holes

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Active Galactic Nuclei



Power supply: gravitational potential energy of accreting material.

Therefore luminosity is the rest mass energy of accreted material multiplied by some efficiency factor: $L = \epsilon \dot{M} c^2$

Eddington Limit

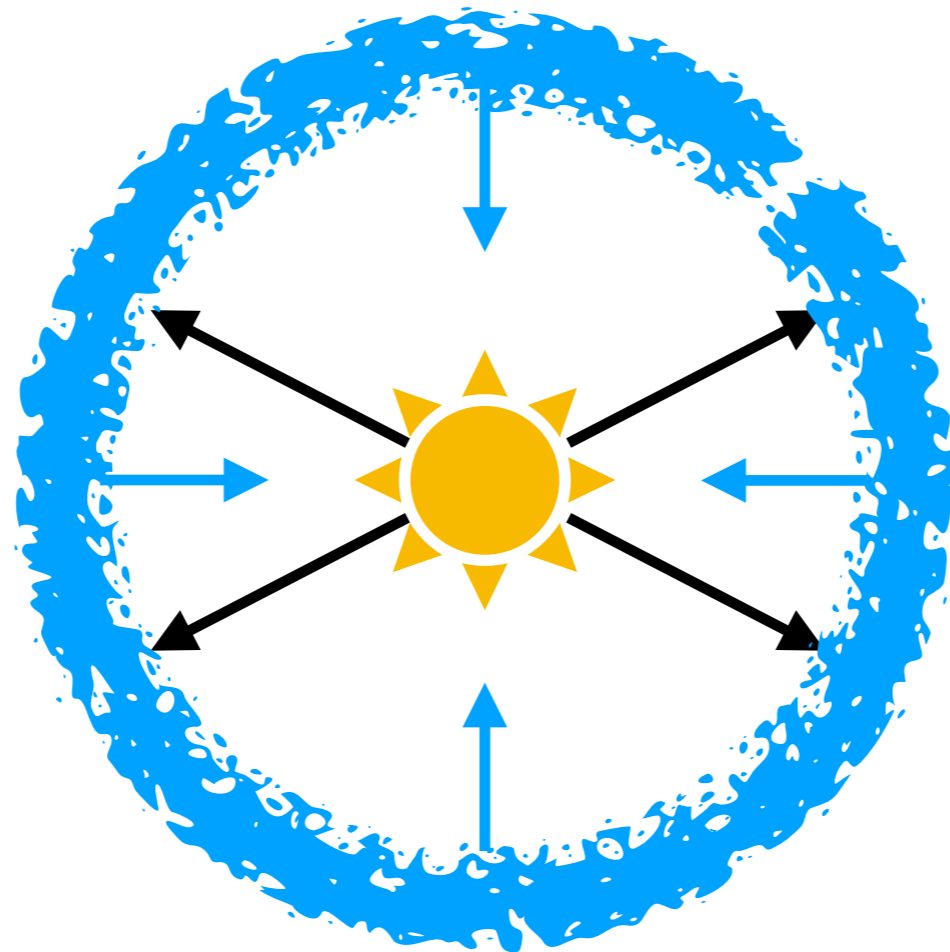
Theoretical maximum luminosity when outwards radiation force balances inwards gravitational force. Above this luminosity, accreting material will be thrown off in winds.

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Assumptions:

- Spherical symmetry
- Dominant opacity = electron scattering (Thomson absorption cross-section).
- Gravity acts predominantly on protons.



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Neutral material therefore can set the two equal for the Eddington limit:

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Therefore can also define Eddington accretion rate:

$$\dot{M}_{\text{Edd}} = \frac{4\pi GMm_p}{\epsilon c \sigma_T}$$

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Therefore AGN can be **much** more luminous than X-ray binaries after all!